

*Domes made of two
dimensional crystals:
magneto- and
quantum-optical properties*

Antonio Polimeni

Dipartimento di Fisica



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UNIVERSITÀ DI ROMA



Acknowledgements I



Salvatore
Cianci

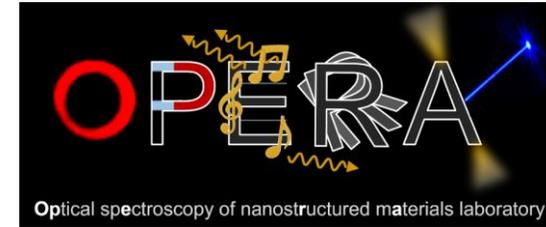
Marzia
Cuccu

Atanu
Patra

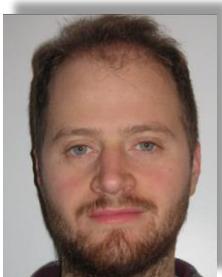
Marco
Felici

Elena
Blundo

Federico
Tuzi



Prof. Antonio Polimeni
Prof. Marco Felici
Elena Blundo
Antonio Miriametro
Marzia Cuccu
Federico Tuzi
Eirini Parmenopoulou
Djeero Peters



CNR IFN
Istituto di Fotonica e Nanotecnologie

Dr. Giorgio Pettinari

Acknowledgements II

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Marco Felici, Marzia Cuccu
Salvatore Cianci, Federico Tuzi
Antonio Miriametro



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National Research Council, Roma, Italy

This project was funded within the QuantERA II Programme that has received funding from the European Union's Horizon 2020 research and innovation programme under Grant Agreement No 101017733"

Acknowledgements II

Elena Blundo

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Salvatore Cianci, Federico Tuzi
Antonio Miriametro



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Cinzia Di Giorgio Fabrizio Bobba



Physics Department,
University of Salerno, Fisciano, Italy

Tanju Yildirim



国立研究開発法人 物質・材料研究機構
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Center for Functional Sensor & Actuator (CFSN), National Institute
for Materials Science (NIMS), Tsukuba, Japan

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Research School of Engineering and College of
Engineering and Computer Science, Canberra, Australia

Katarzyna Olkowska-Pucko Tomasz Kazimierczuk Adam Babiński Maciej Molas

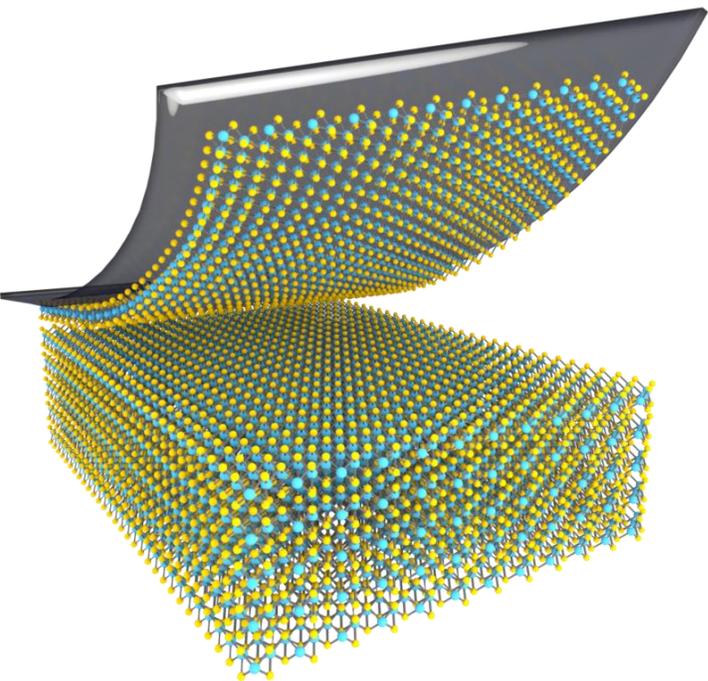


Institute of Experimental Physics, Faculty of Physics,
University of Warsaw, Poland

Outline

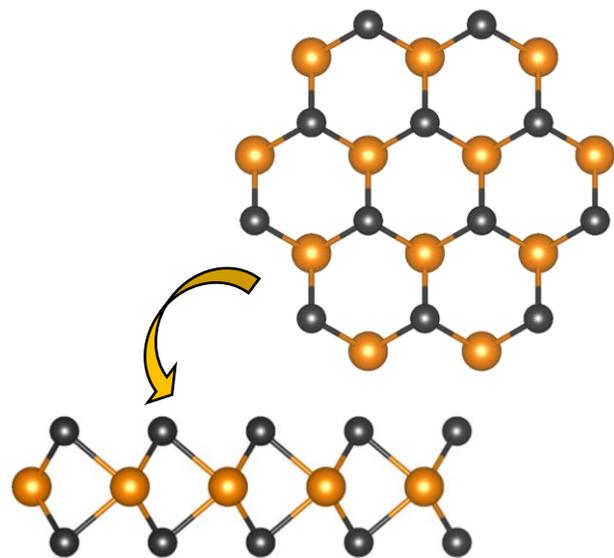
- Strain in 2D crystals
- Formation, characteristics and control of artificial domes in exfoliable materials
- Strain fields in curved membranes: optical, vibrational and magneto-optical properties
- Applications for site-controlled quantum light sources

Two-dimensional crystals

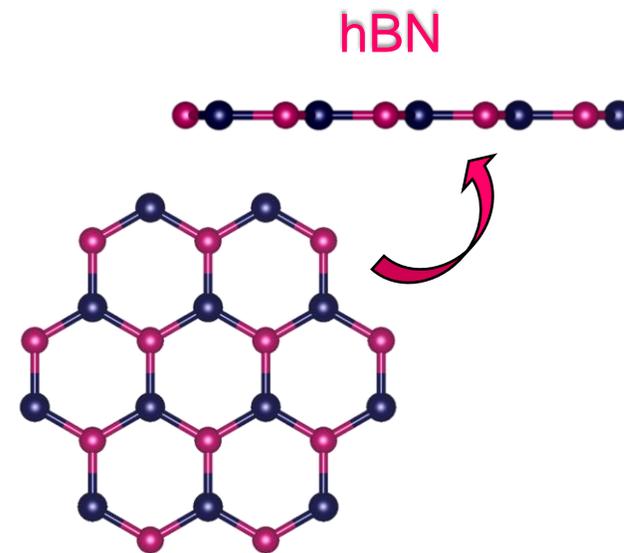


M: Mo, W

X: S, Se, Te

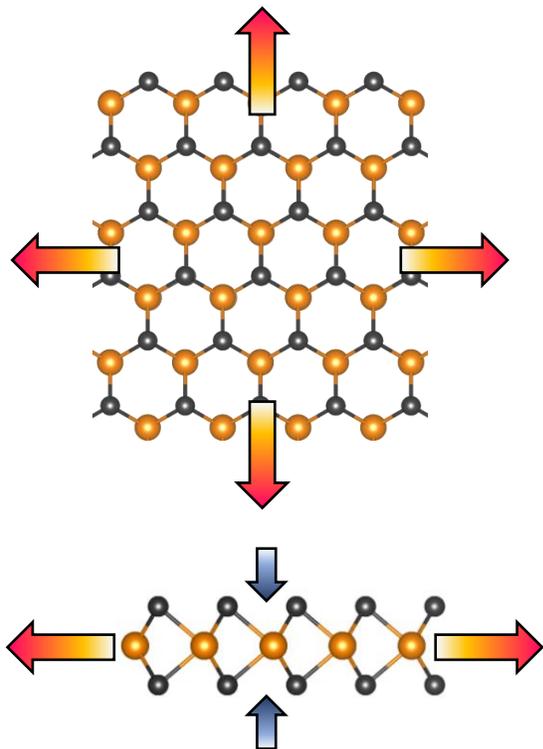


TMDs

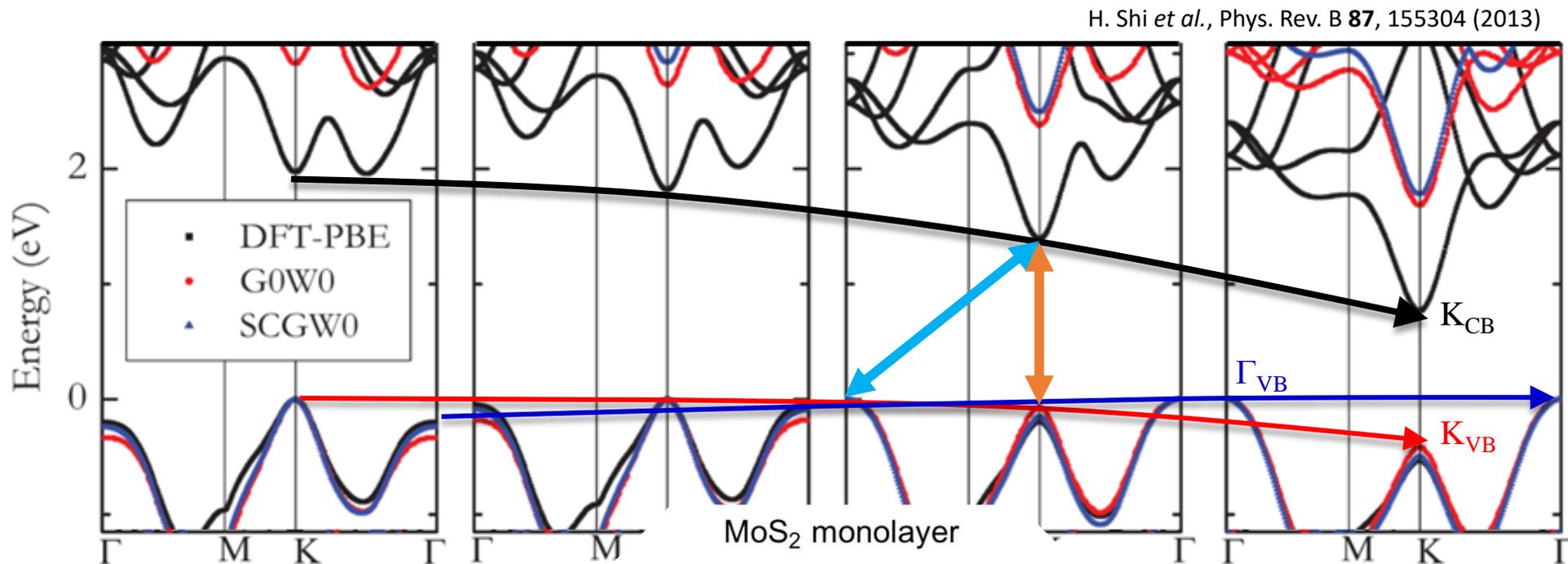


Why strain?

WS₂ monolayer

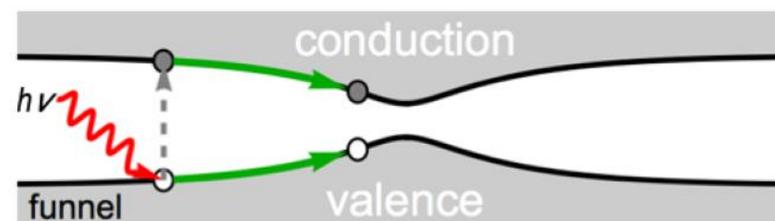
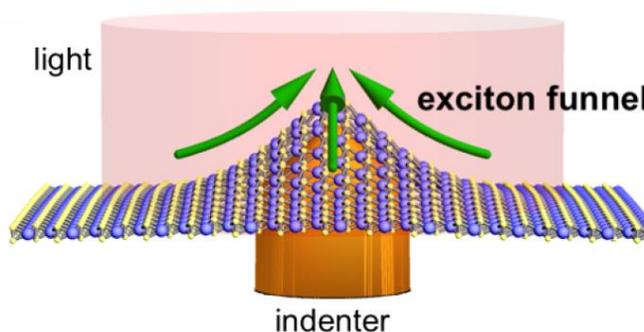


ε in-plane strain

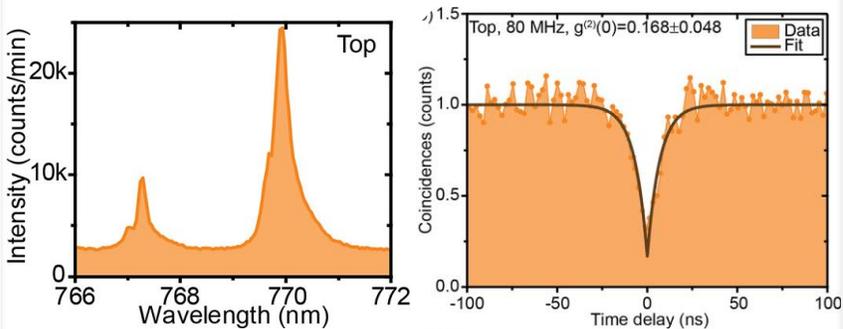
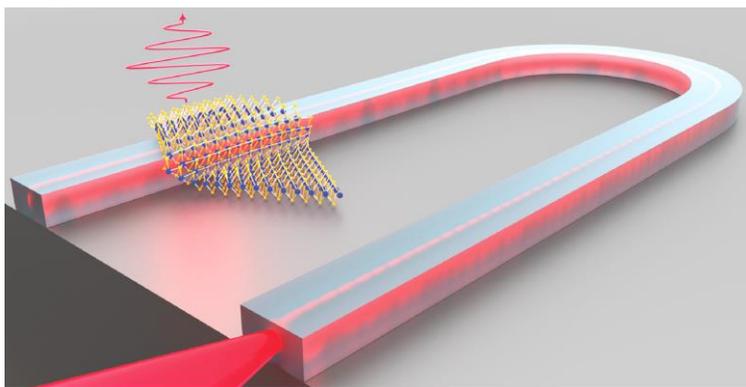


$\varepsilon = 0\%$

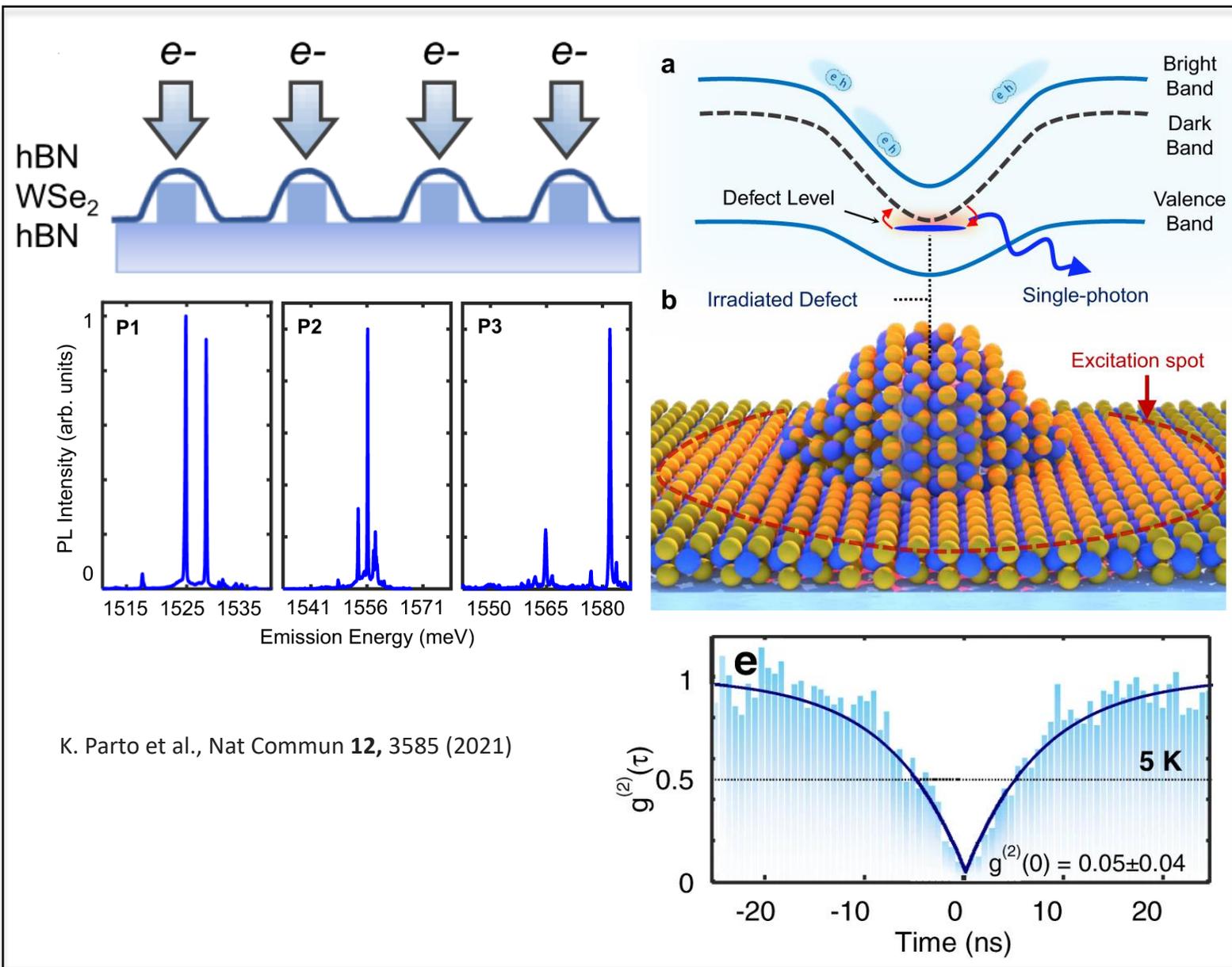
$\varepsilon = 12\%$



Why strain?



Carlos Errando-Herranz et al., ACS Photonics **8**, 1069 (2021)

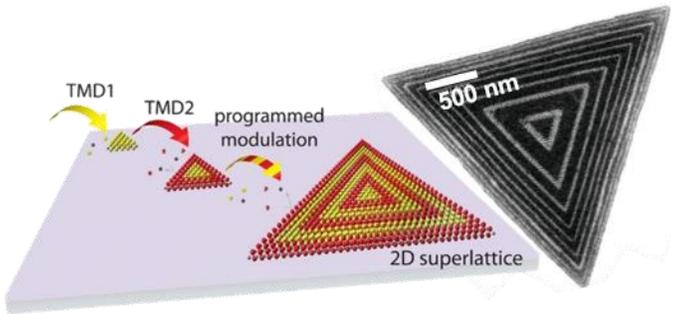


K. Parto et al., Nat Commun **12**, 3585 (2021)

Straining methods

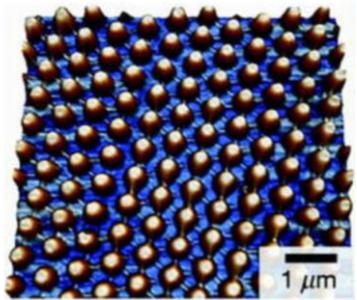
E. Blundo et al.
Appl. Phys. Rev. **8**, 021318 (2021)

Growth of superlattices

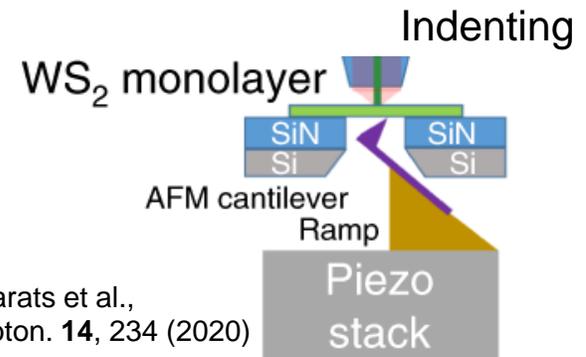


Xie et al., Science **359**, 1131 (2020)

Deposition on nanocones/pillars

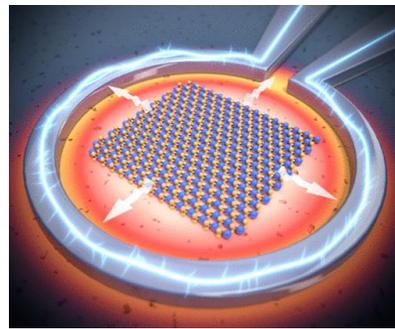


Li et al., J. Am. Chem. Soc. **138**, 5123 (2016)



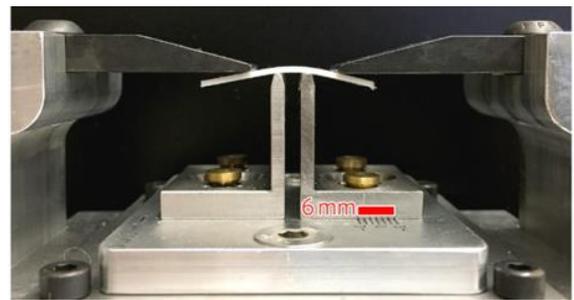
M. G. Harats et al.,
Nat. Photon. **14**, 234 (2020)

Stretching



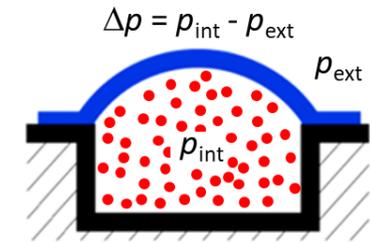
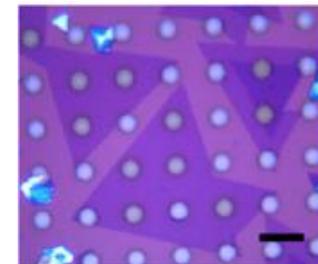
Ryu et al., Nano Lett. **20**, 5339 (2020)

Bending



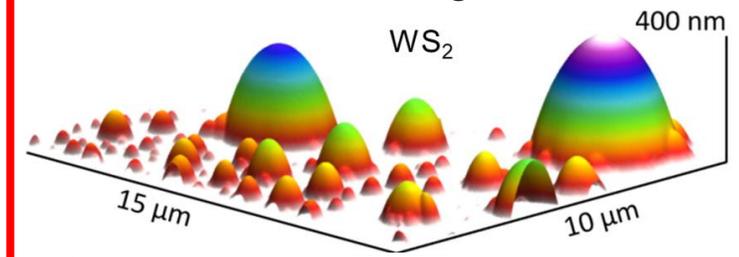
Wu et al., Nano Lett. **18**, 2351 (2018)

Bulging



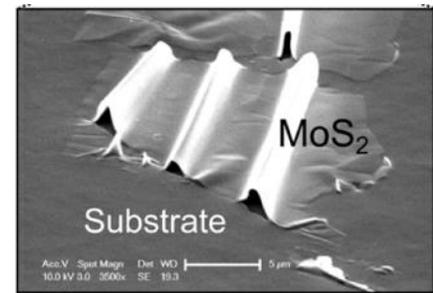
Lloyd et al., Nano Lett. **16**, 5836 (2016)

Bubbling



Blundo et al, Phys. Rev. Research **2**, 012024 (2020)

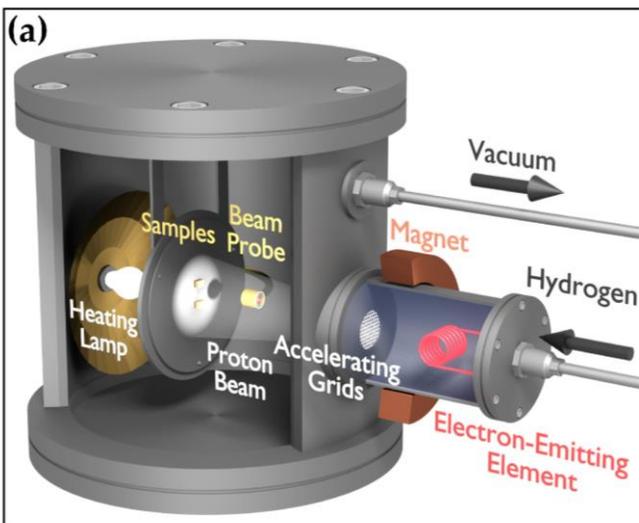
Wrinkling



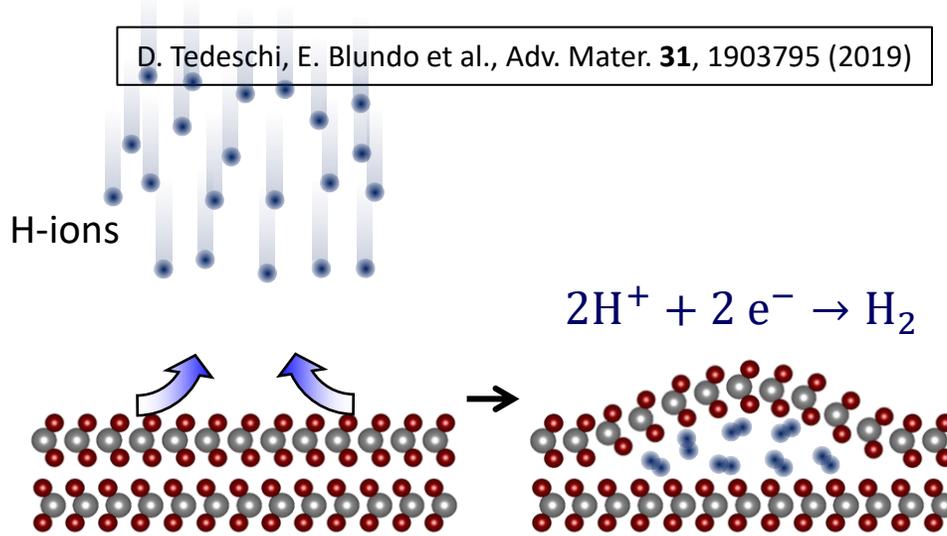
Castellanos-Gomez et al., Nano Lett. **13**, 5361 (2013)

Hydrogen irradiation of bulk TMDs

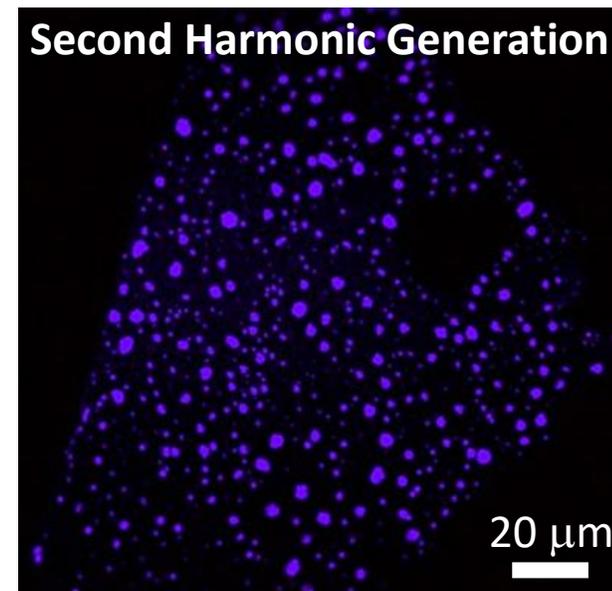
D. Tedeschi, E. Blundo et al., Adv. Mater. **31**, 1903795 (2019)



J. Felton, E. Blundo et al., Molecules **25**, 2526 (2020)

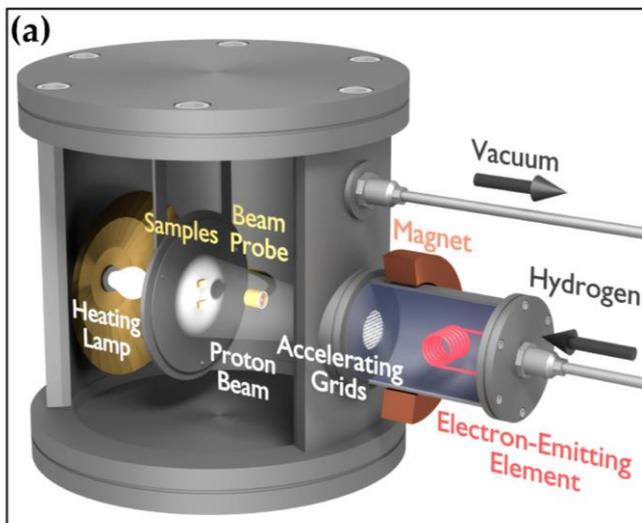


1 monolayer thick !

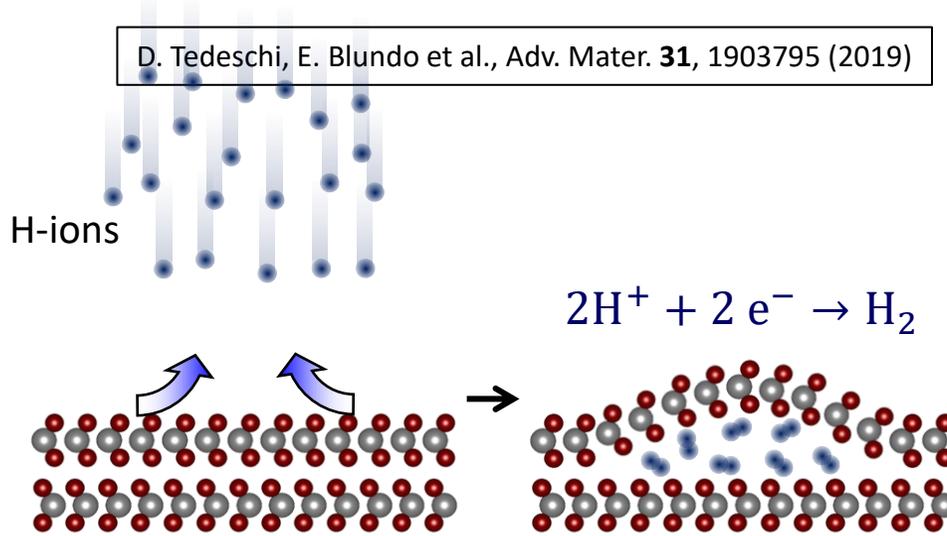


Hydrogen irradiation of bulk TMDs

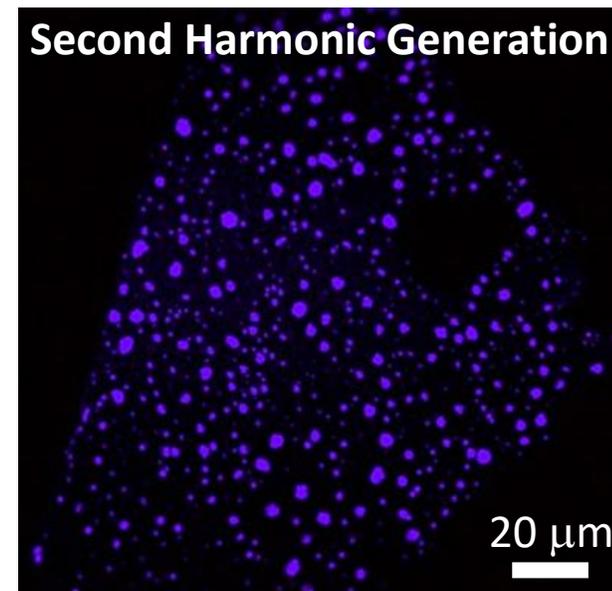
D. Tedeschi, E. Blundo et al., Adv. Mater. **31**, 1903795 (2019)



J. Felton, E. Blundo et al., Molecules **25**, 2526 (2020)



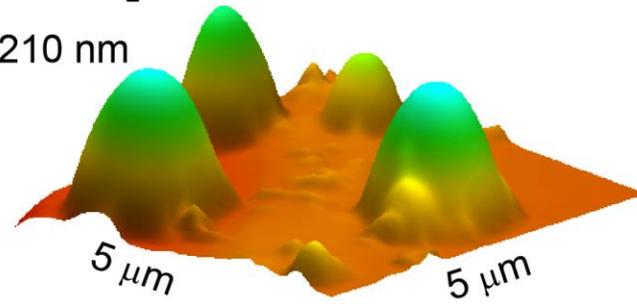
1 monolayer thick !



Domes in TMDs

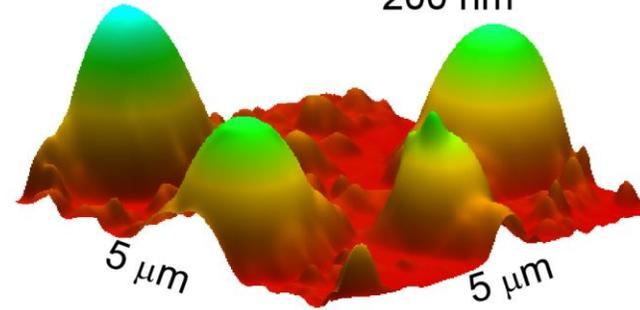
A WS_2

210 nm



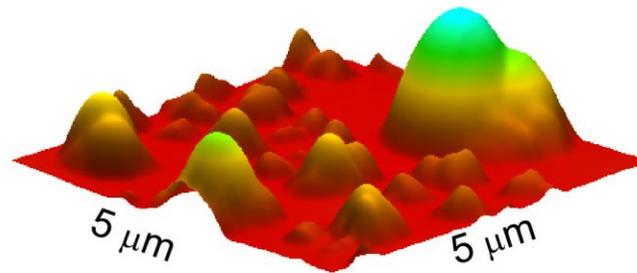
B MoS_2

200 nm



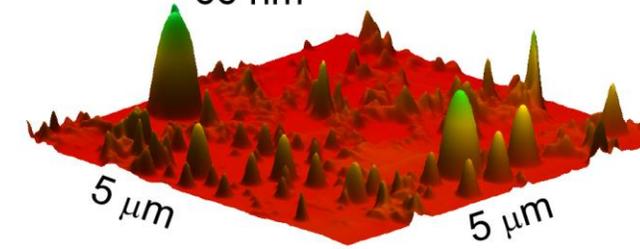
C WSe_2

150 nm



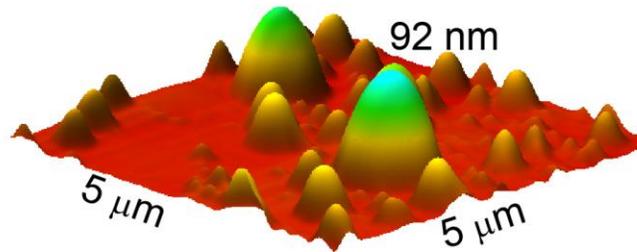
D MoSe_2

66 nm



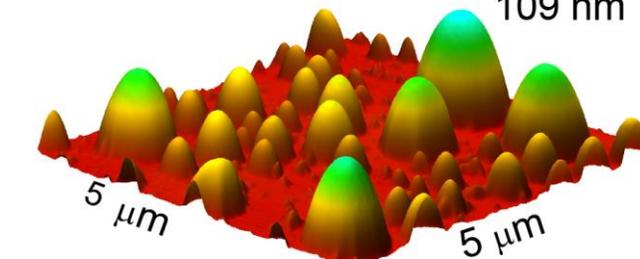
E WTe_2

92 nm

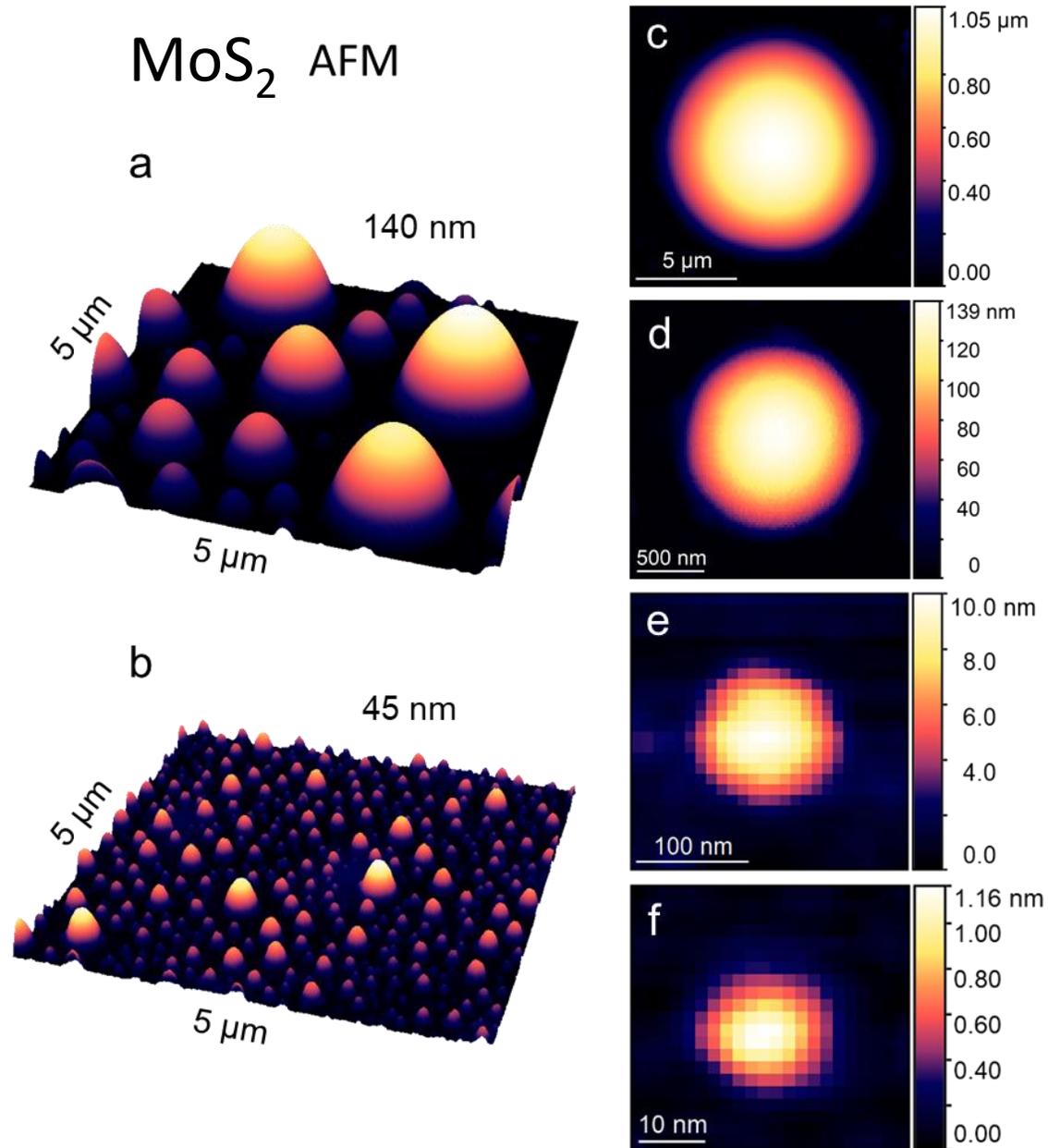


F MoTe_2

109 nm



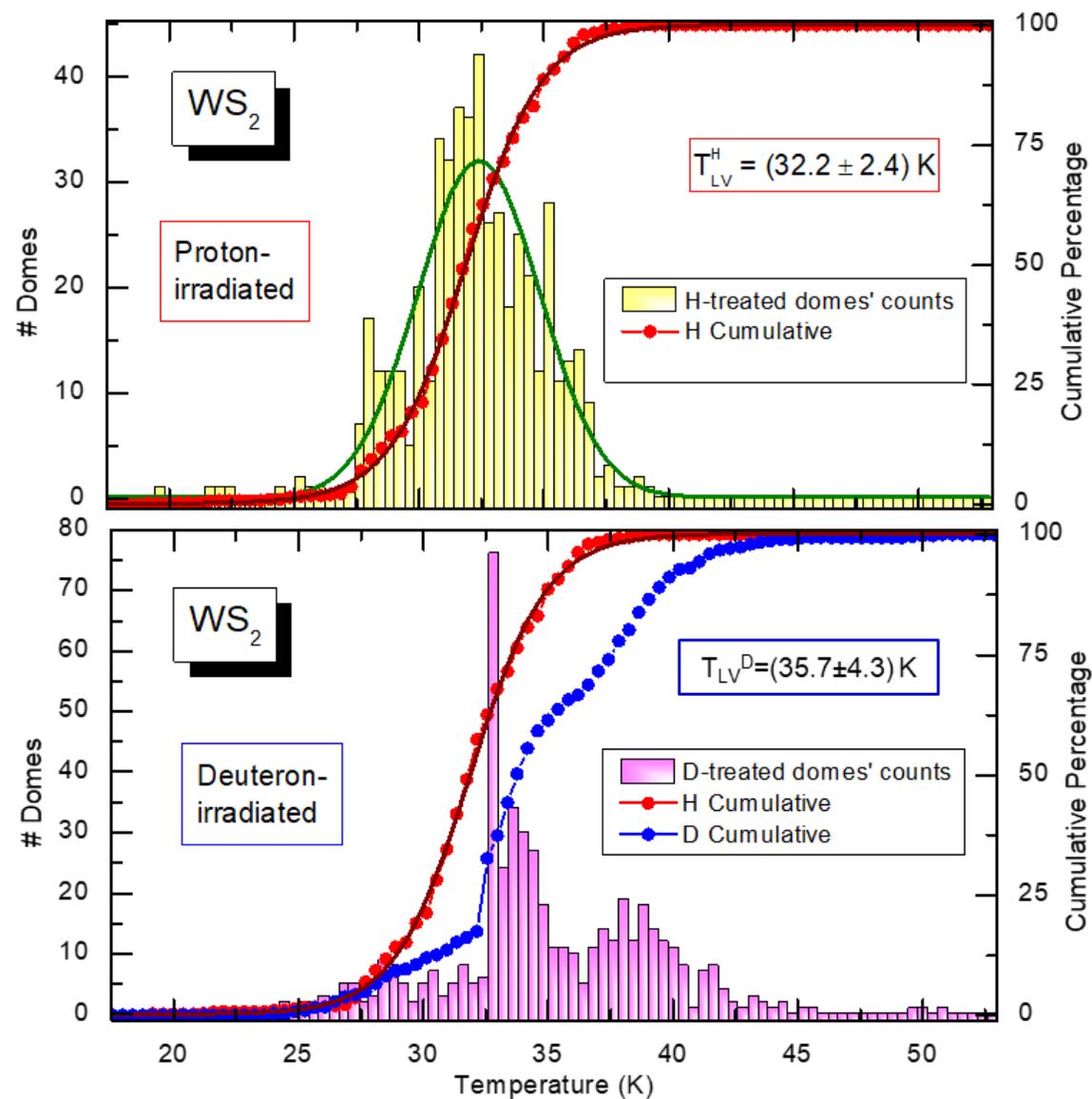
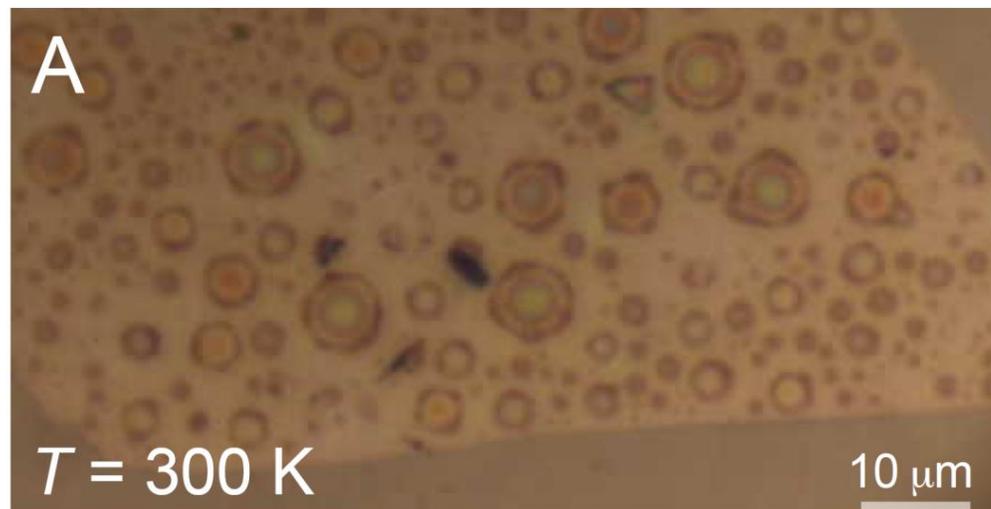
Domes in TMDs: Size



Three orders of magnitude variation in size is possible

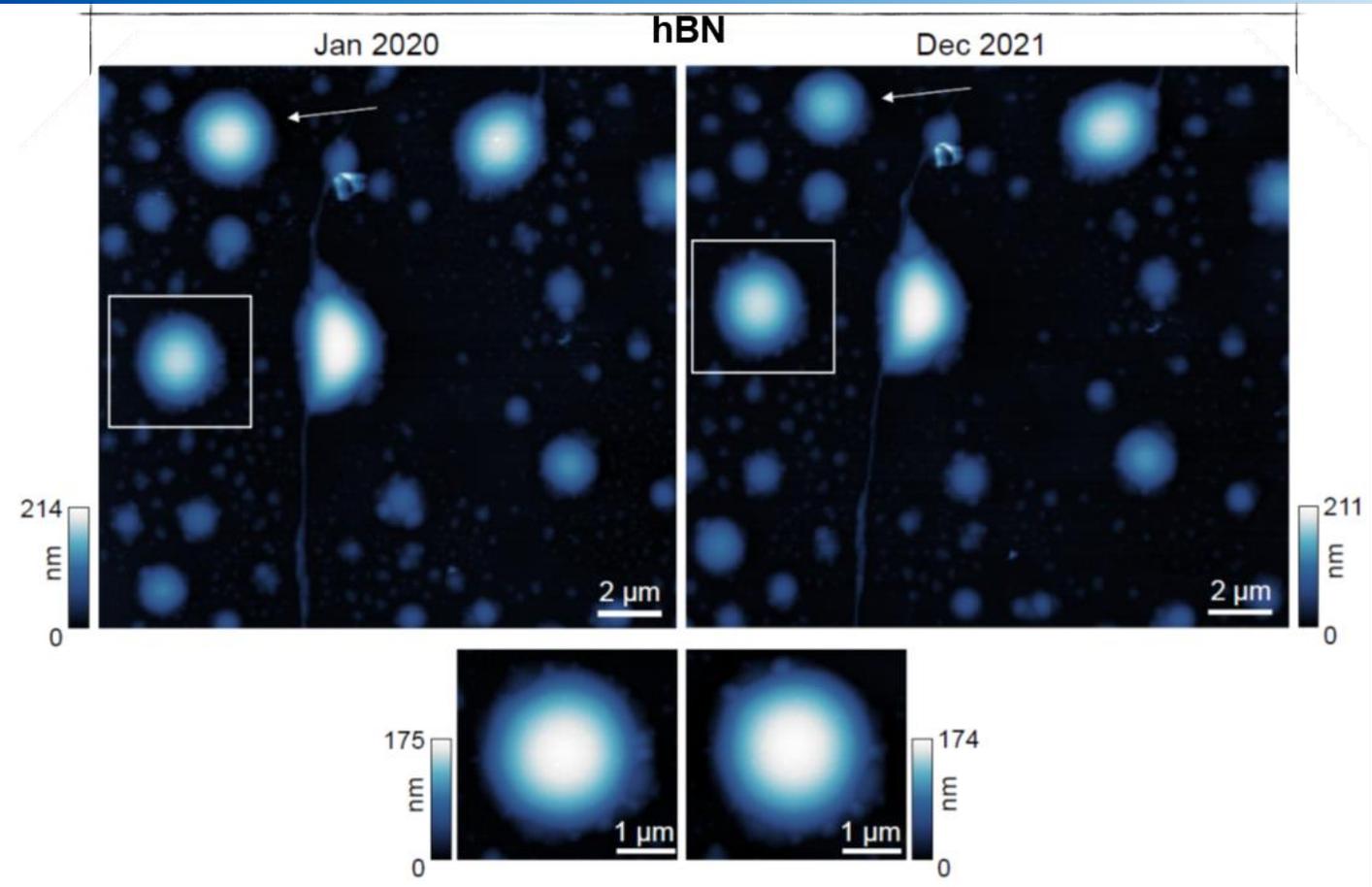
Evidence of H_2 molecules within the domes

$$d_H = 8 \times 10^{16} \text{ protons/cm}^2$$



Domes in TMDs: Main features

Durability > 3 yrs

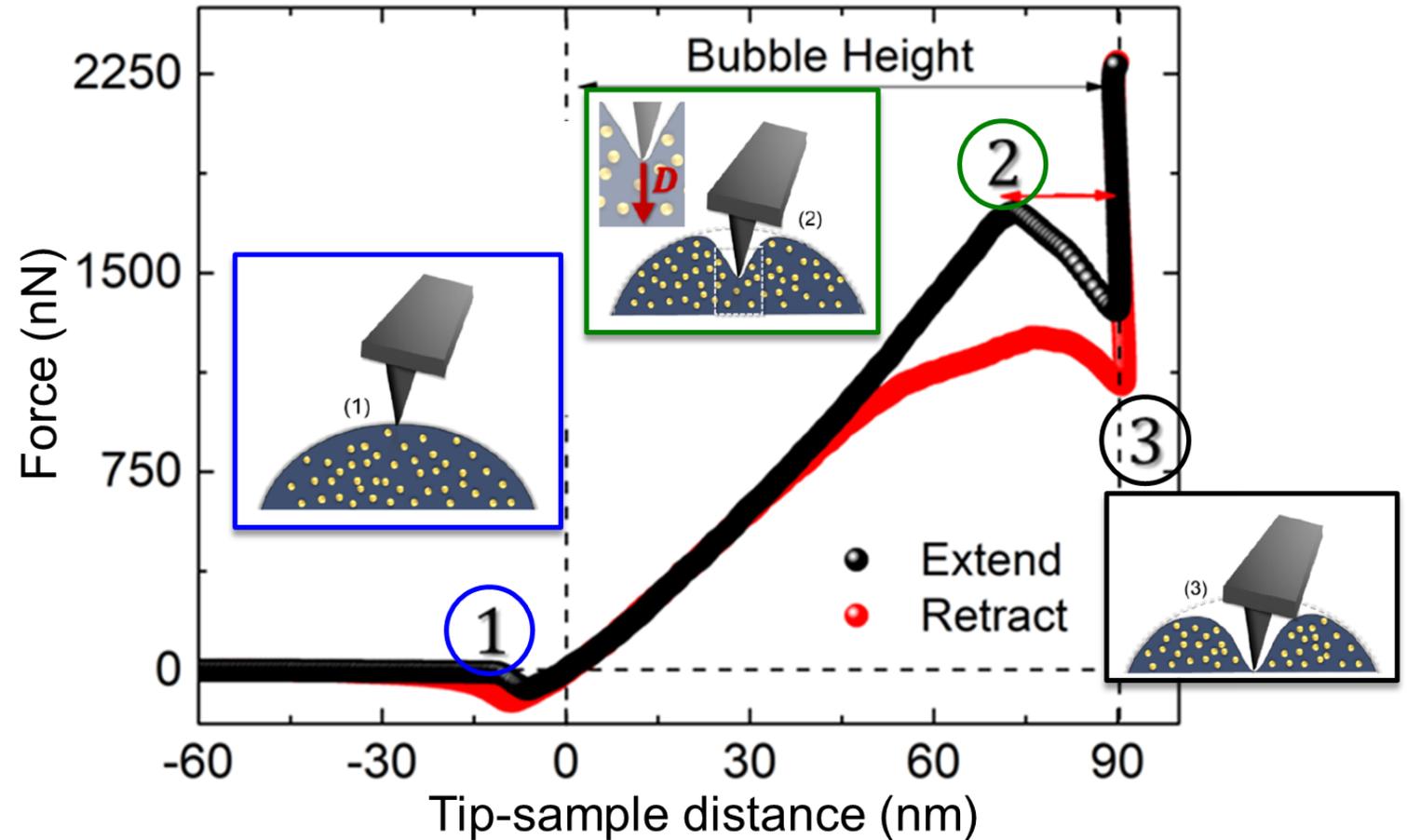
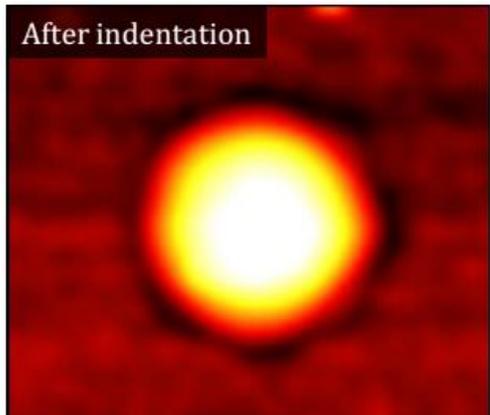
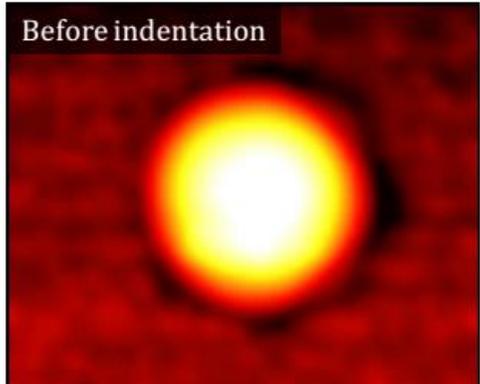


E. Blundo et al., Nano Lett. **22**, 1525 (2022)

Domes in TMDs: Main features

Durability > 3 yrs

Robustness



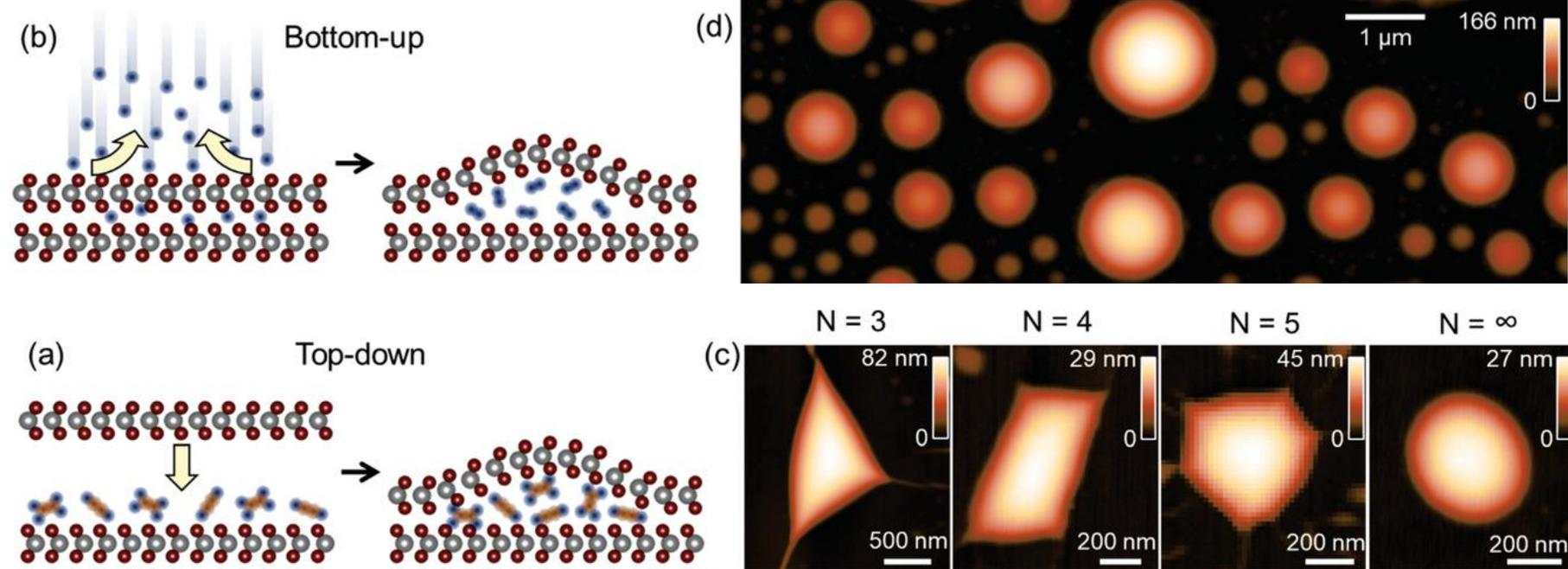
C. Di Giorgio et al., ACS Appl. Mater. Interfaces **13**, 48228 (2021)

Domes in TMDs: Main features

Durability > 3 yrs

Robustness

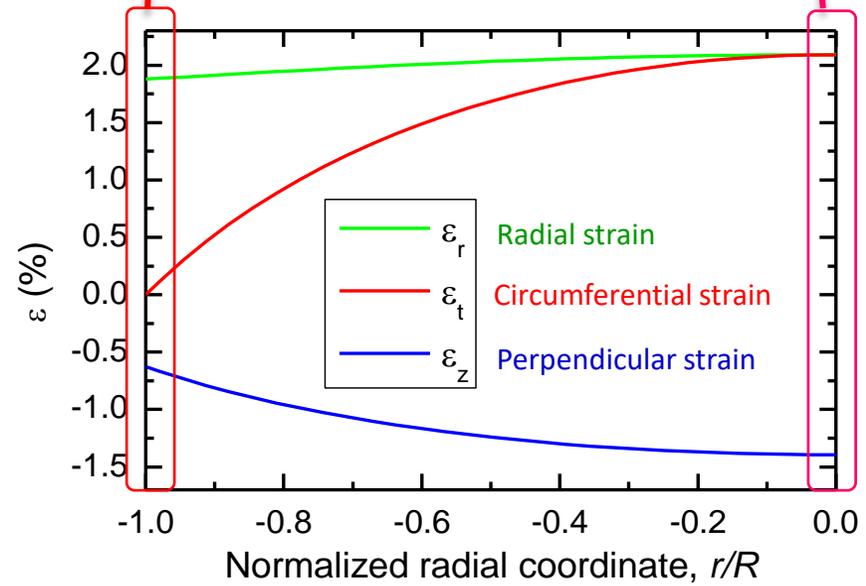
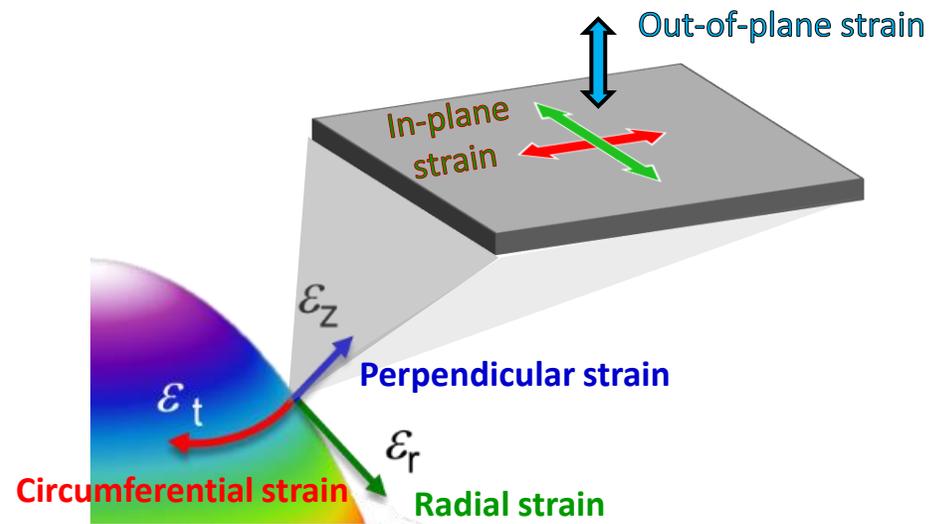
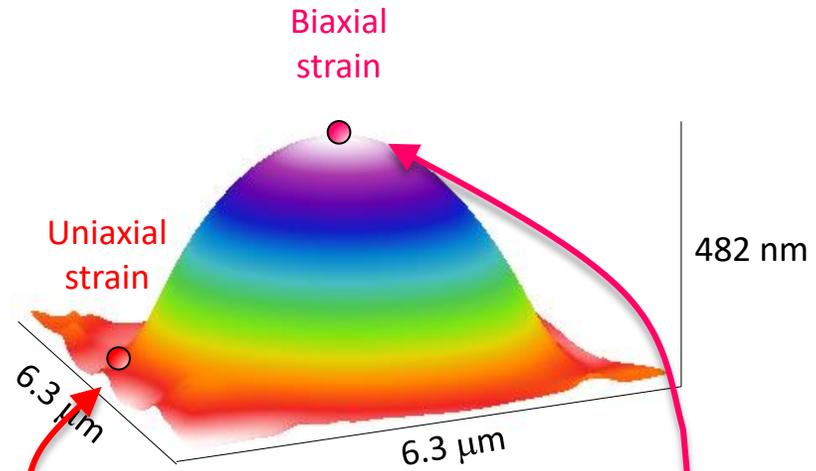
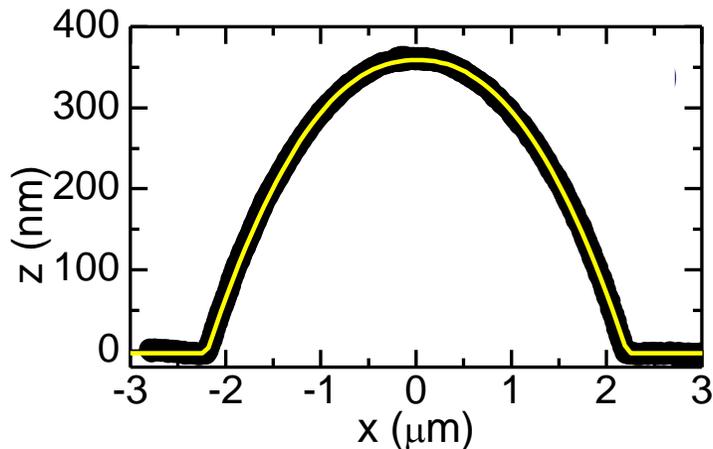
Regular/
reproducible shape



Strain field

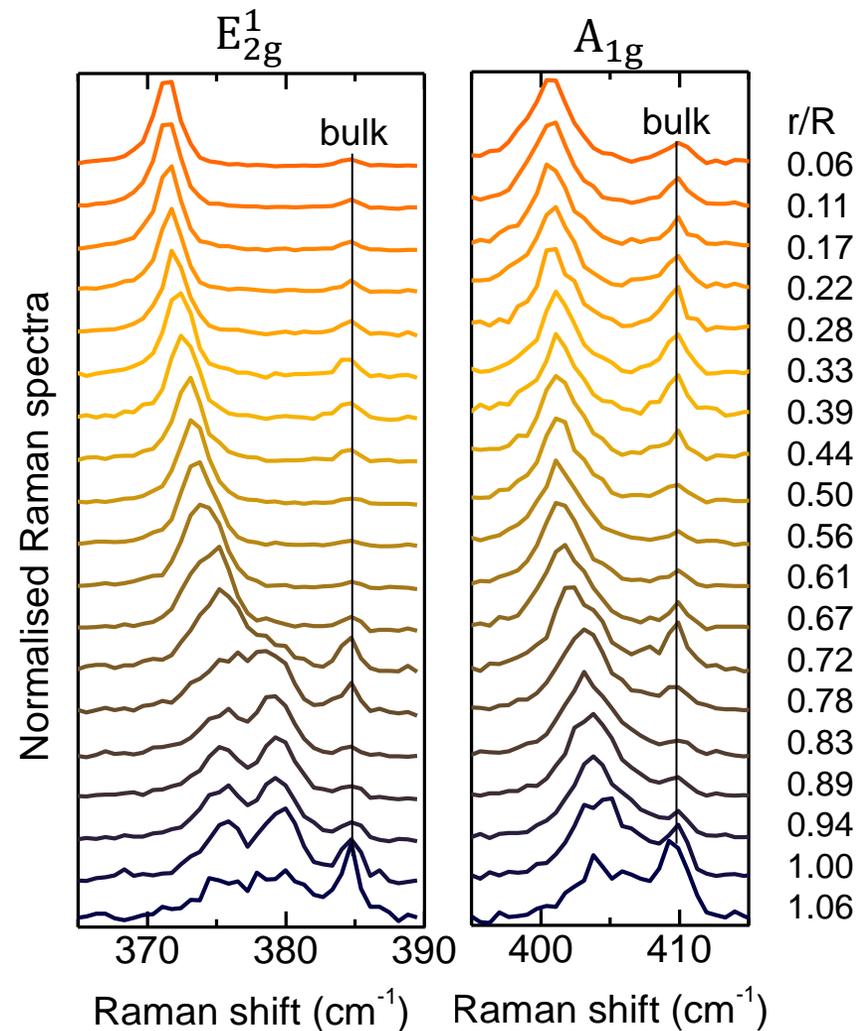
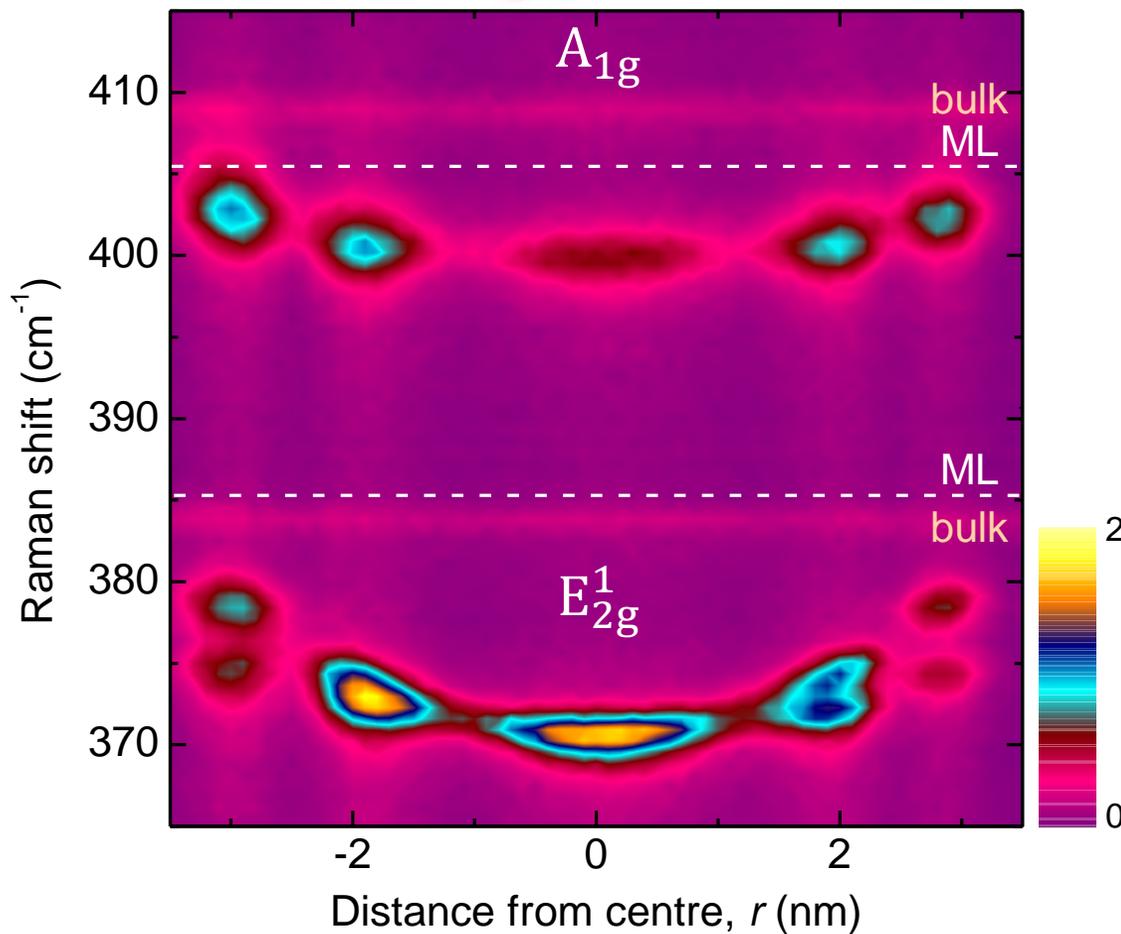
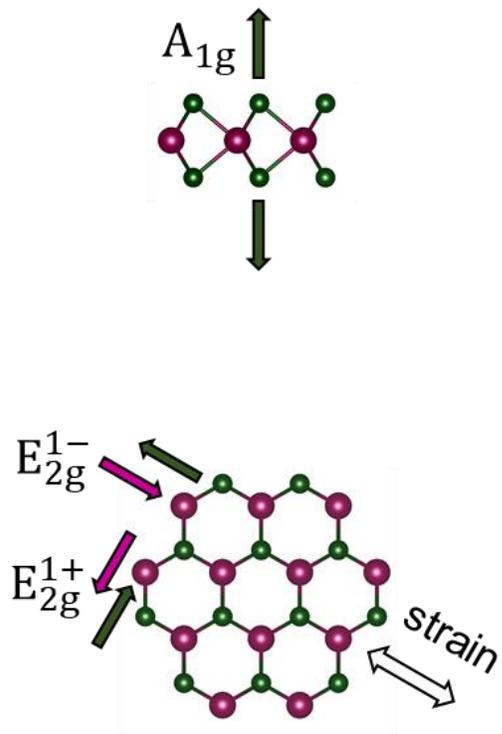
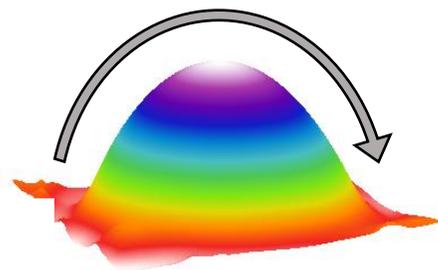
FEM (Finite-Element Method) calculations, within the framework of nonlinear membrane theory

E. Blundo et al, Phys. Rev. Lett. **127**, 46101 (2021)



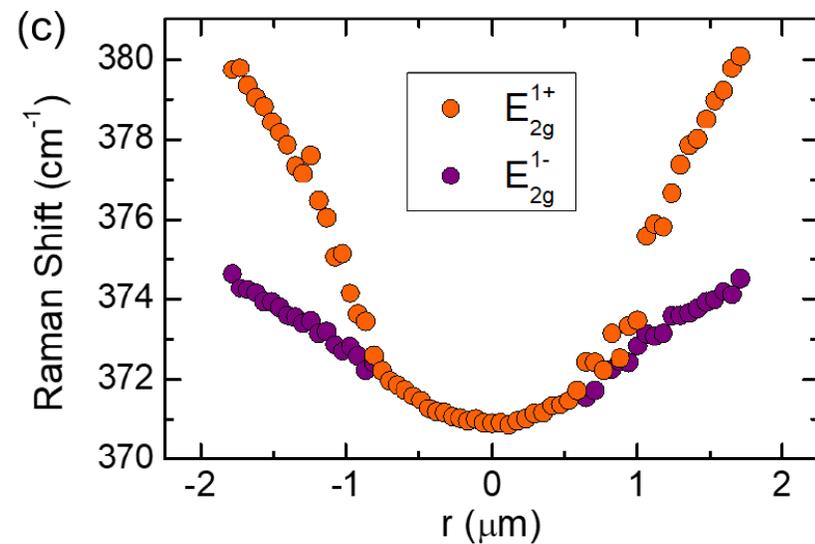
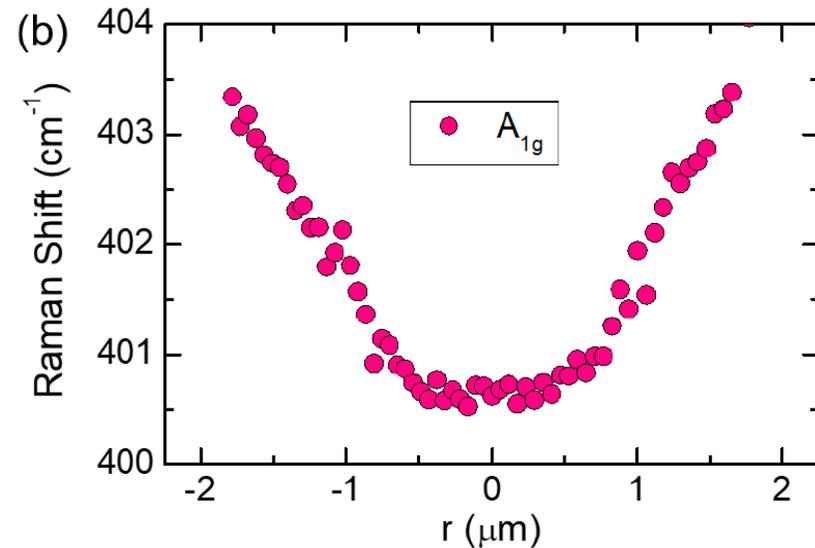
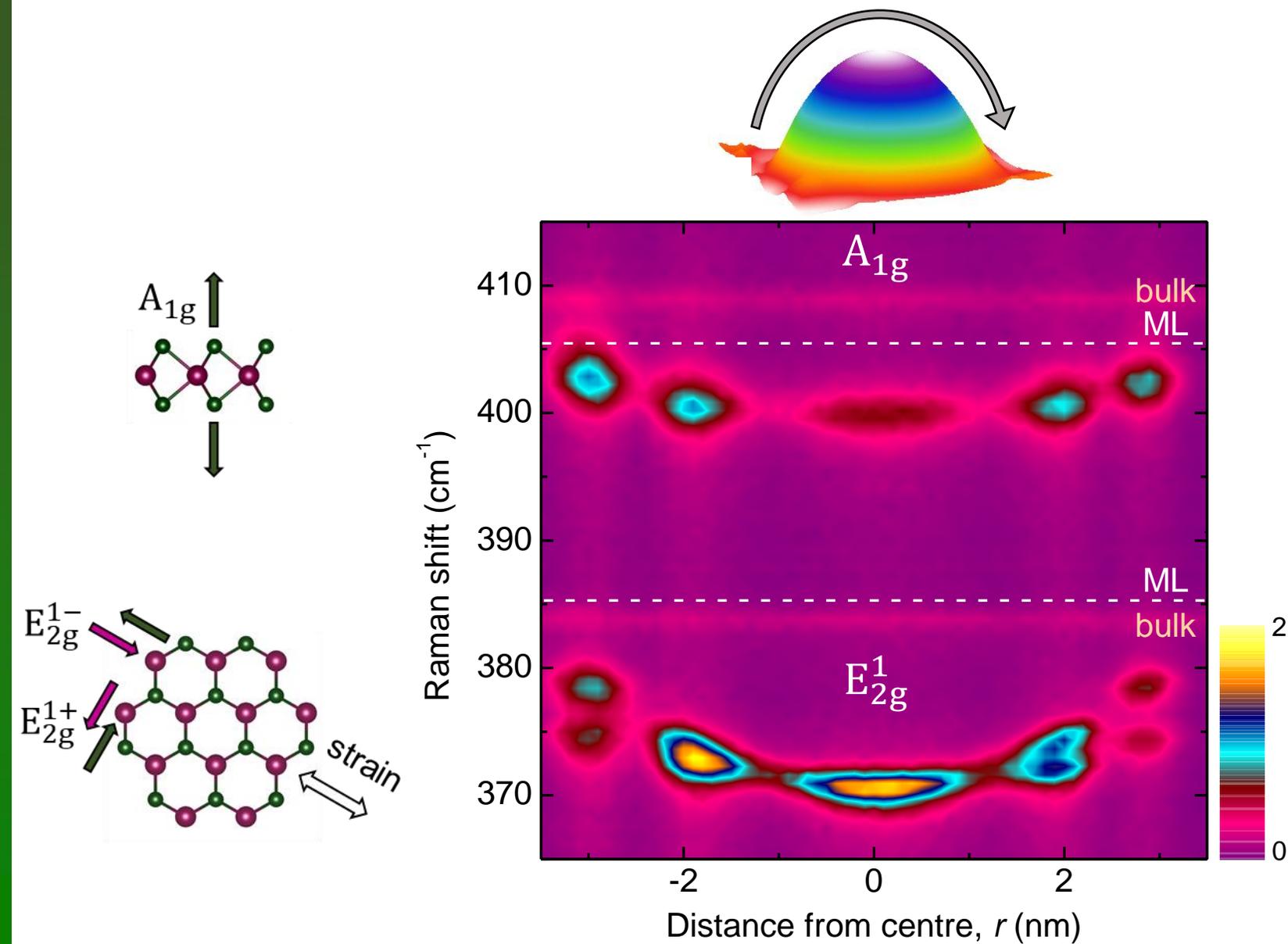
Micro-Raman measurements – MoS₂

E. Blundo et al, Phys. Rev. Lett. **127**, 46101 (2021)

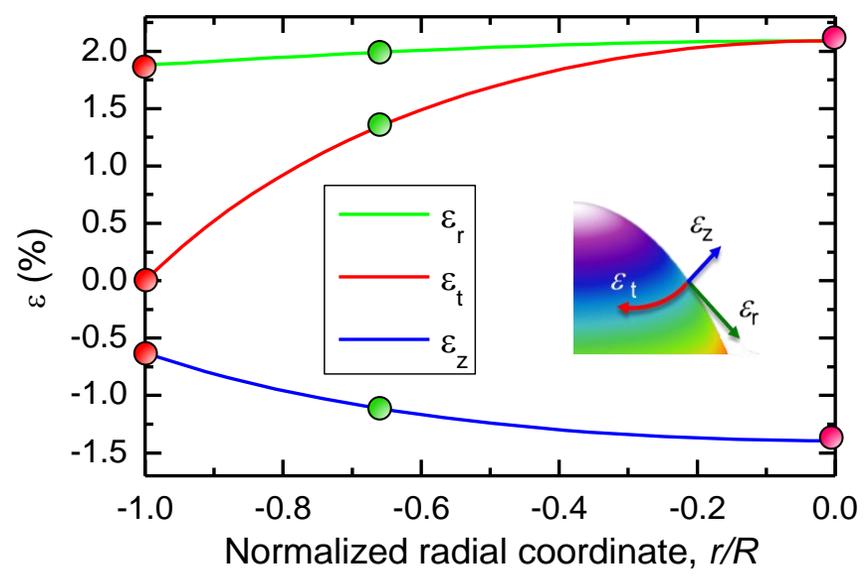
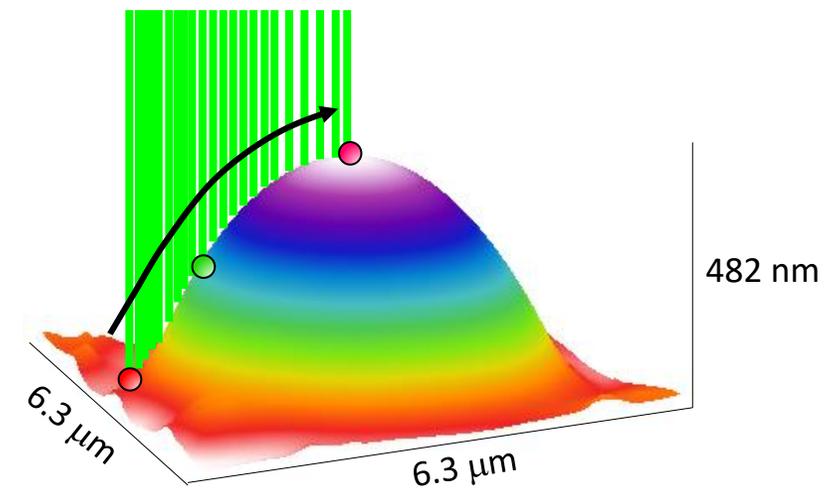
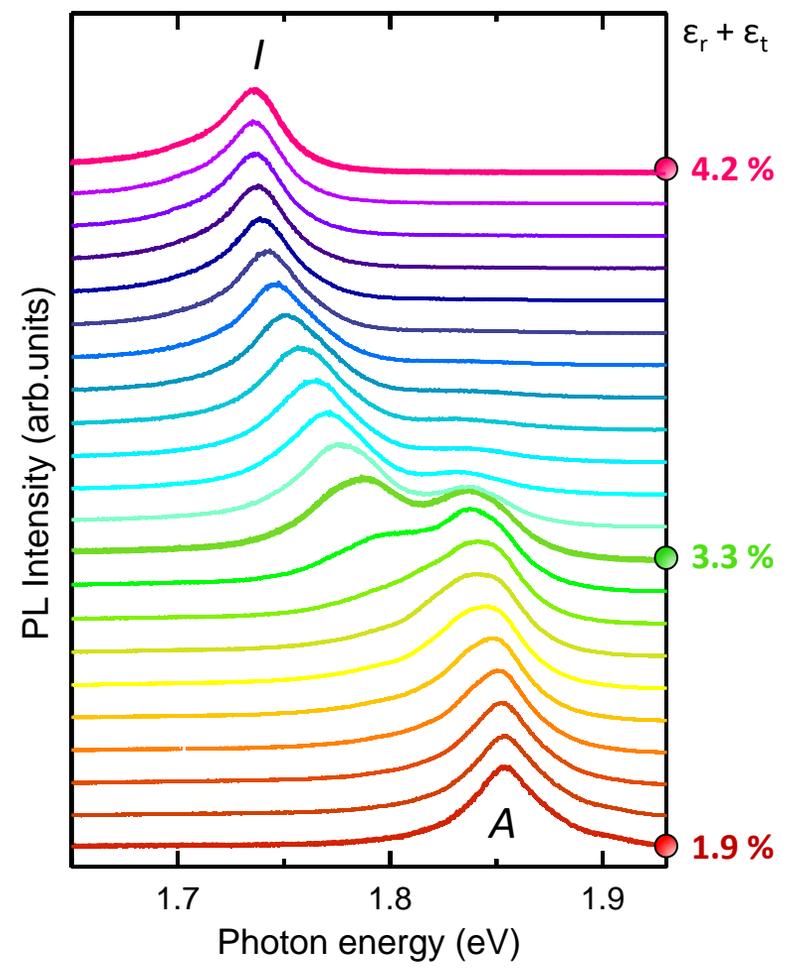


Micro-Raman measurements – MoS₂

E. Blundo et al, Phys. Rev. Lett. **127**, 46101 (2021)

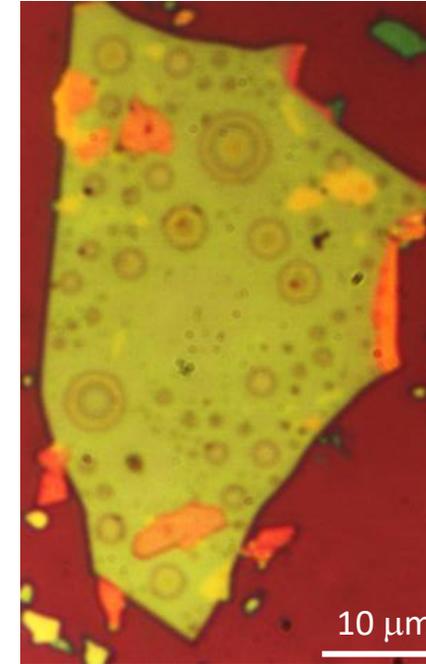
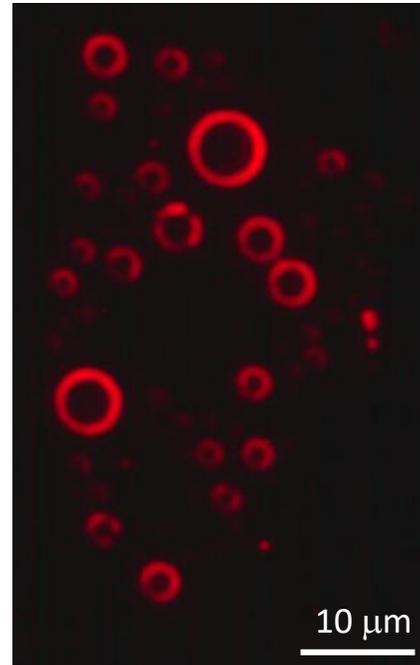
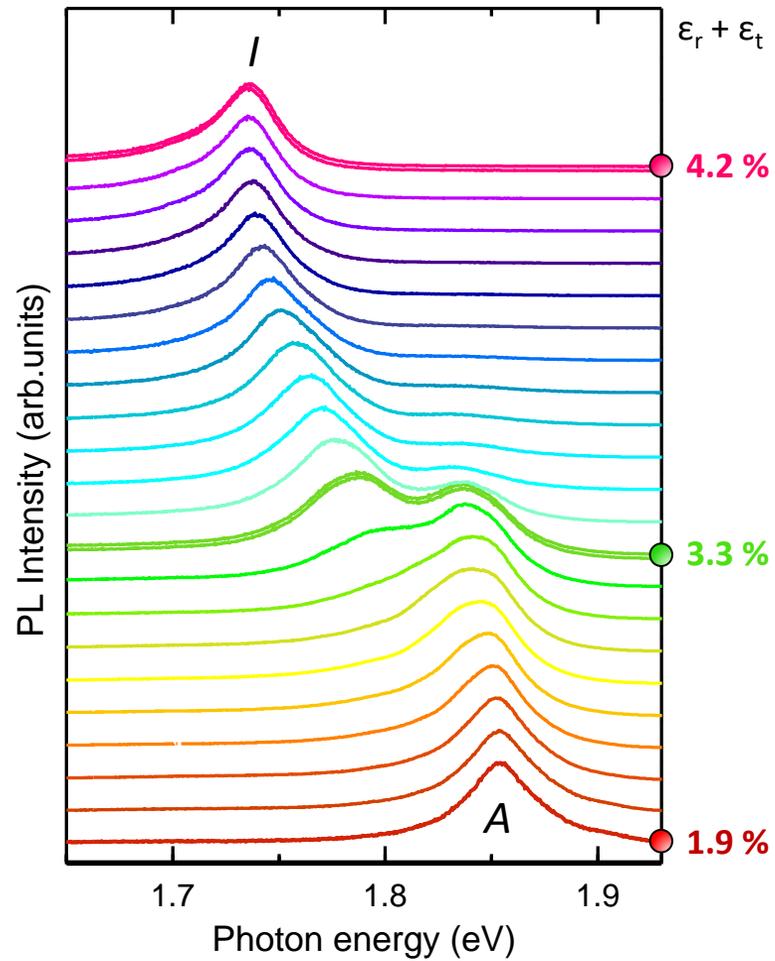


Micro-PL mapping of WS_2 domes



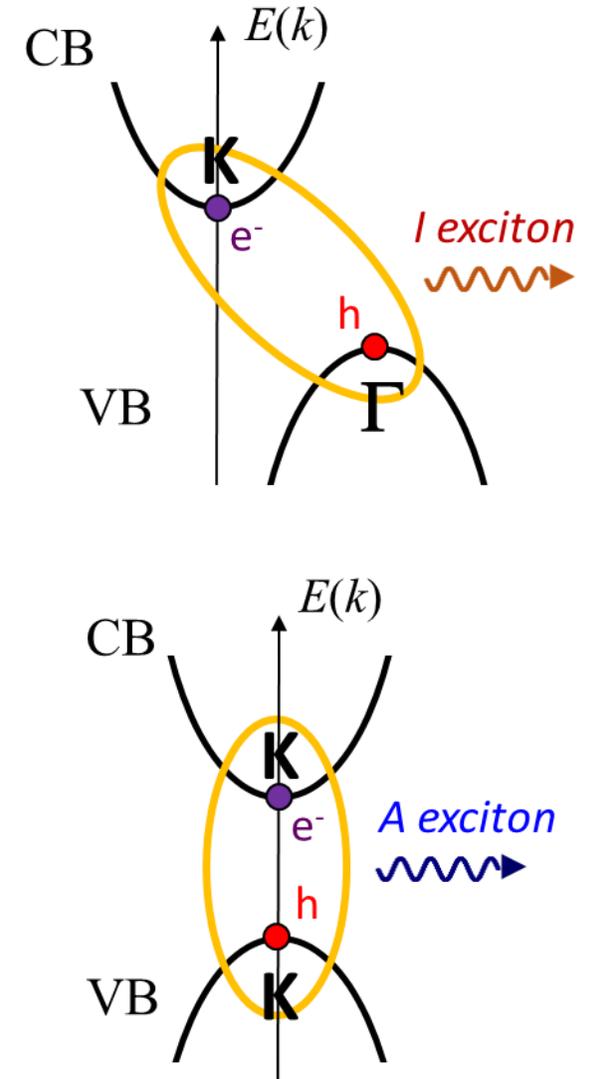
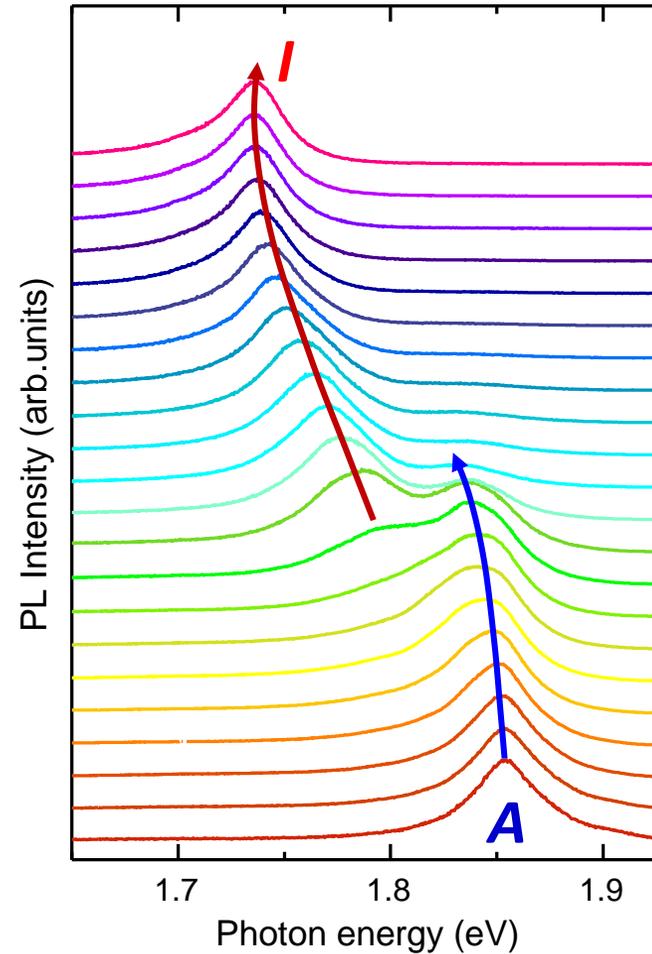
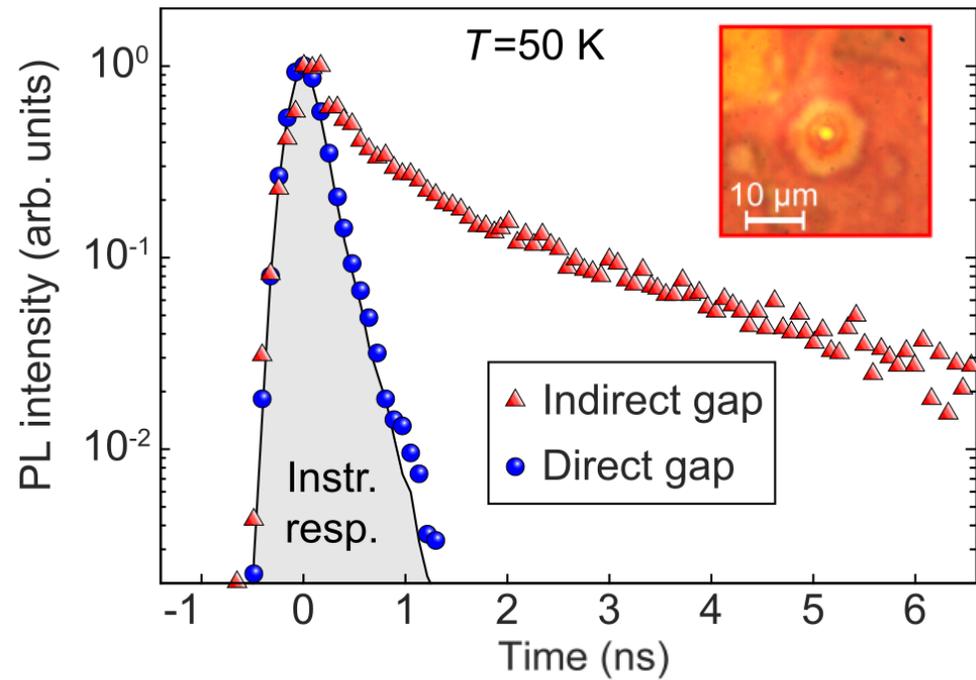
E. Blundo et al., Phys. Rev. Res. 2, 012024 (2020)

Micro-PL mapping of WS_2 domes



Creation of highly efficient anular regions

Micro-PL mapping of WS_2 domes



Magneto-PL measurements of WS_2 domes

Proposal for magnet time at the
High Field Magnet Laboratory
(up to 30 T). PI: A. Polimeni



Proposal for magnet time (**ISABEL**) at
the **Regional Partner 'University of
Warsaw'** (up to 12 T). PI: E. Blundo

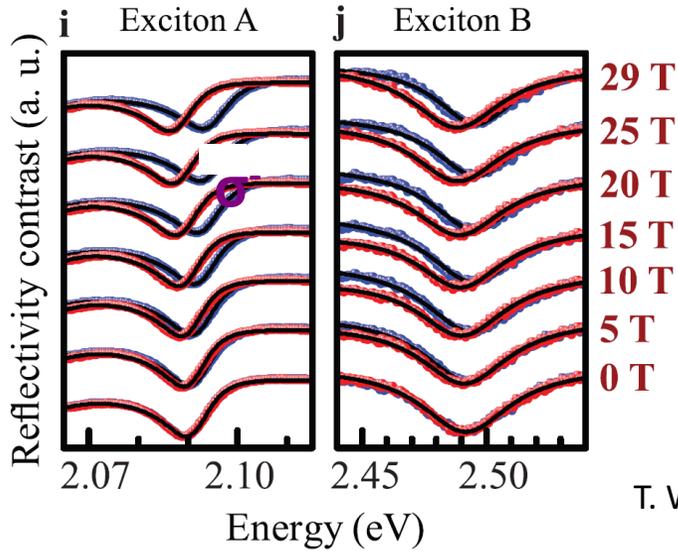


Magneto-PL measurements of WS_2 monolayers

WS_2 monolayer

M. Koperski et al., 2D Mater. **6**, 015001 (2019)

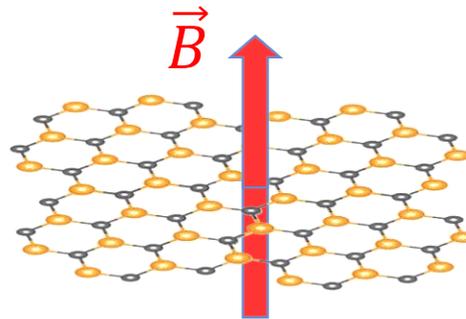
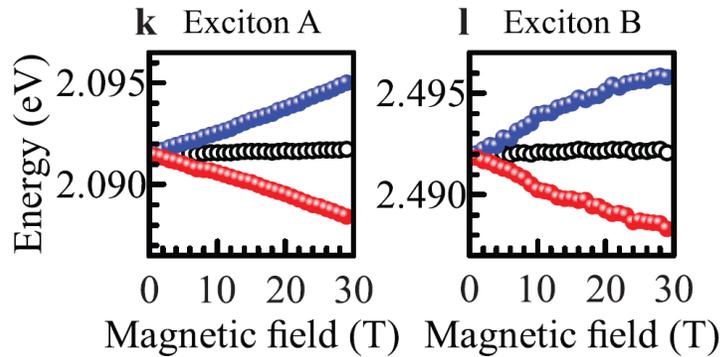
E. Blundo et al., Phys. Rev. Lett. **129**, 067402 (2022)



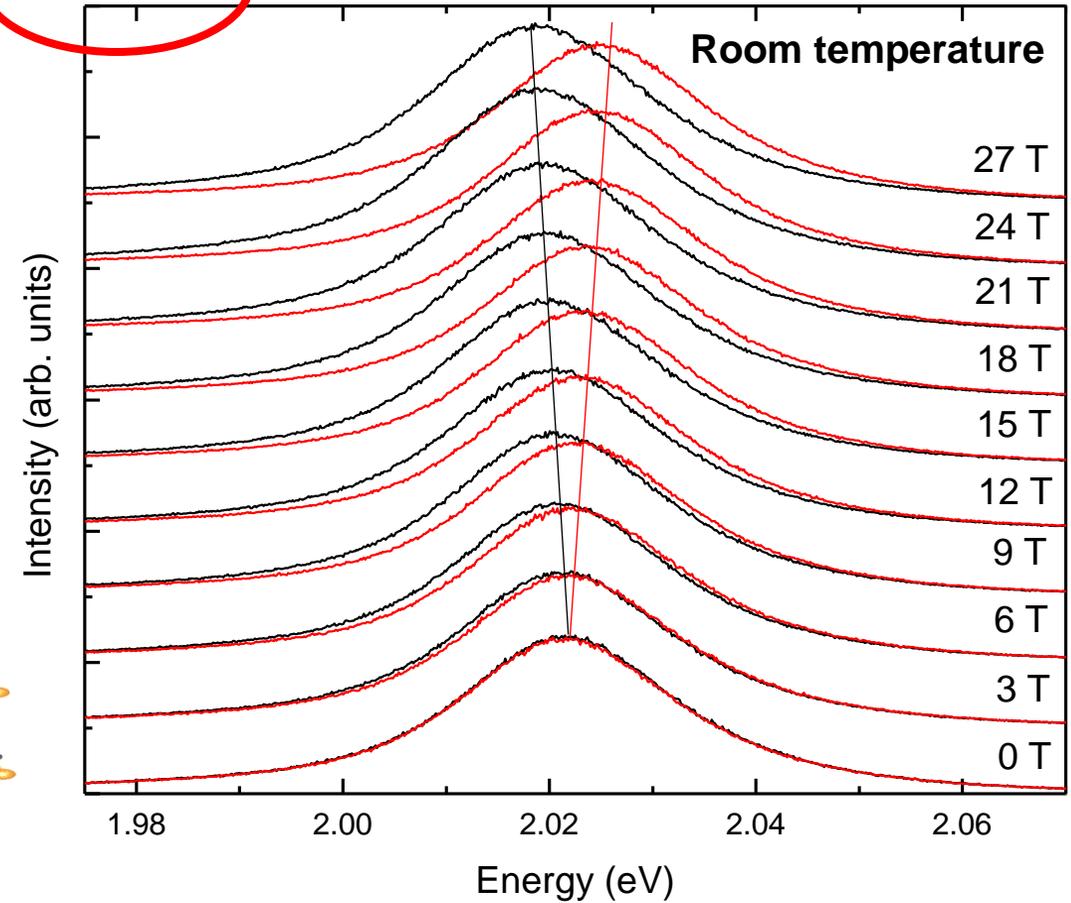
$$\Delta E_Z^A = g_A \mu_B B$$

$$g_A = 2[L_Z(CB_+) - L_Z(VB_+)]$$

T. Woźniak et al., Phys. Rev. B **101**, 235408 (2020)



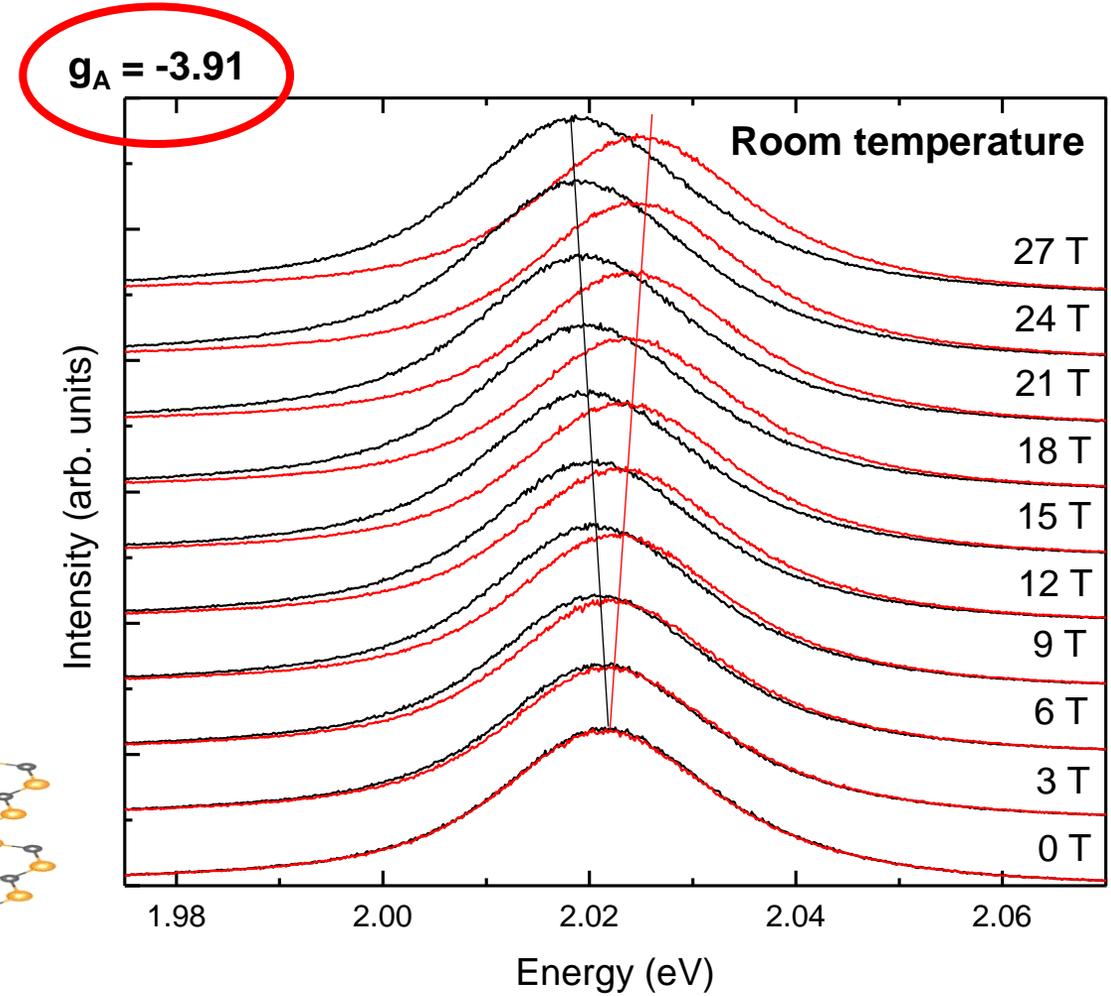
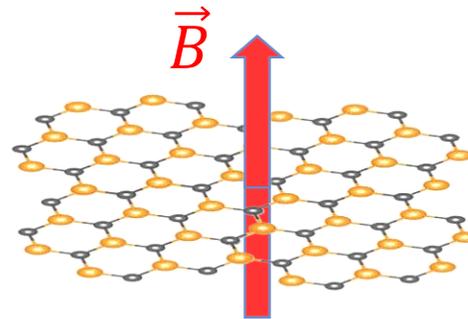
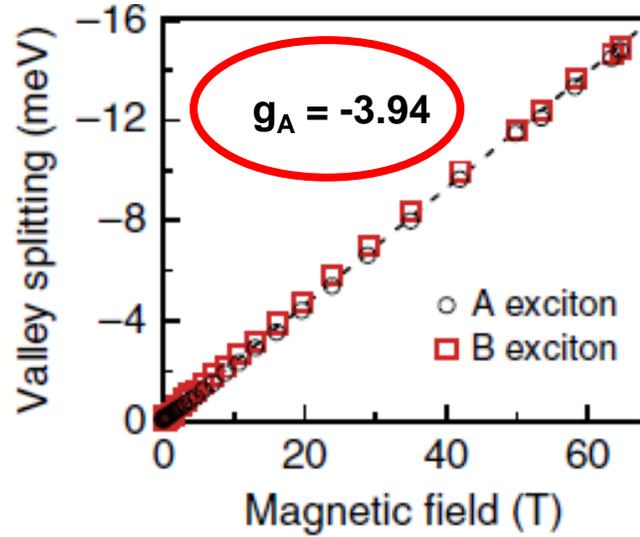
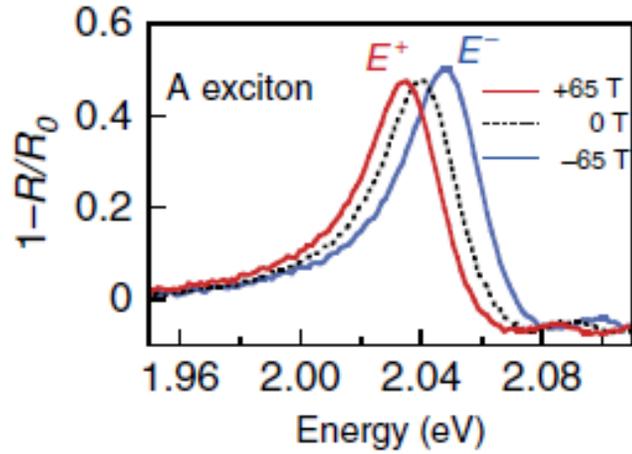
$g_A = -3.91$



Line	g-factor
Exciton A	-3.7 ± 0.2
Exciton B	-4.9 ± 1.2

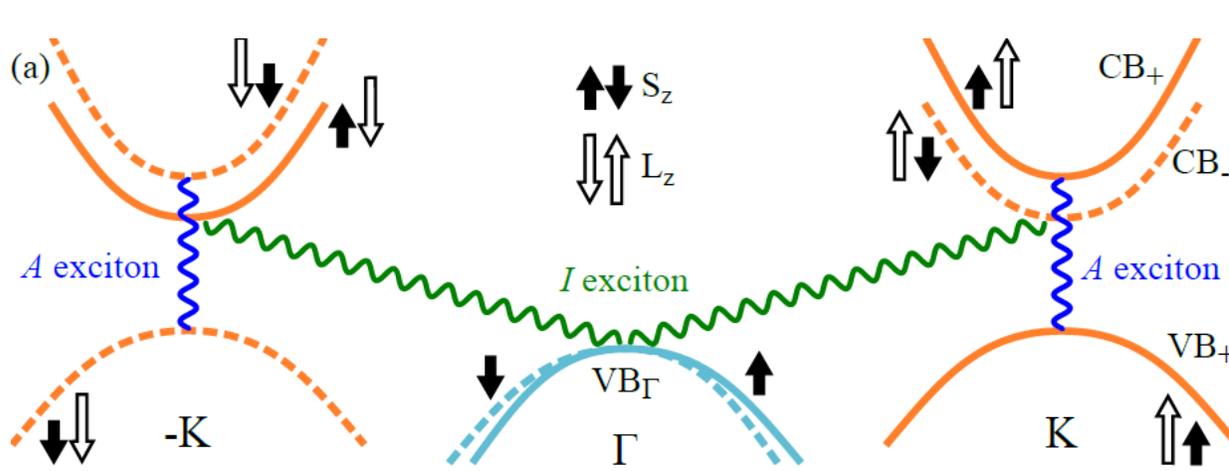
Magneto-PL measurements of WS_2 monolayers

E. Blundo et al., Phys. Rev. Lett. **129**, 067402 (2022)



Stier et al., Nat. Commun. **7**, 10643 (2016)

Magneto-PL measurements of WS_2 domes

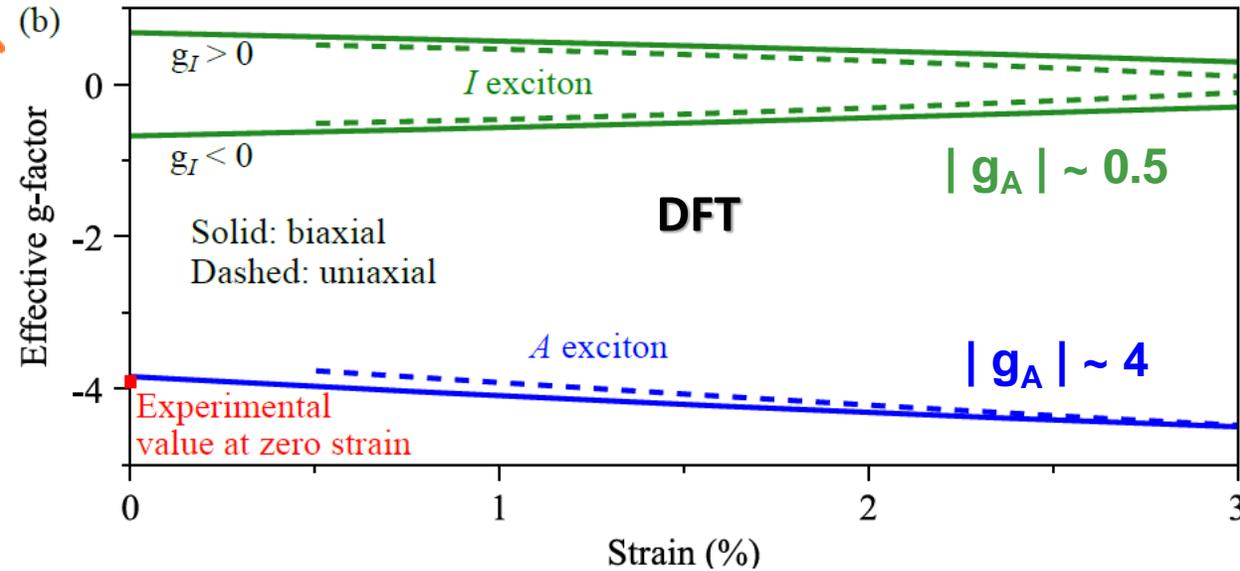
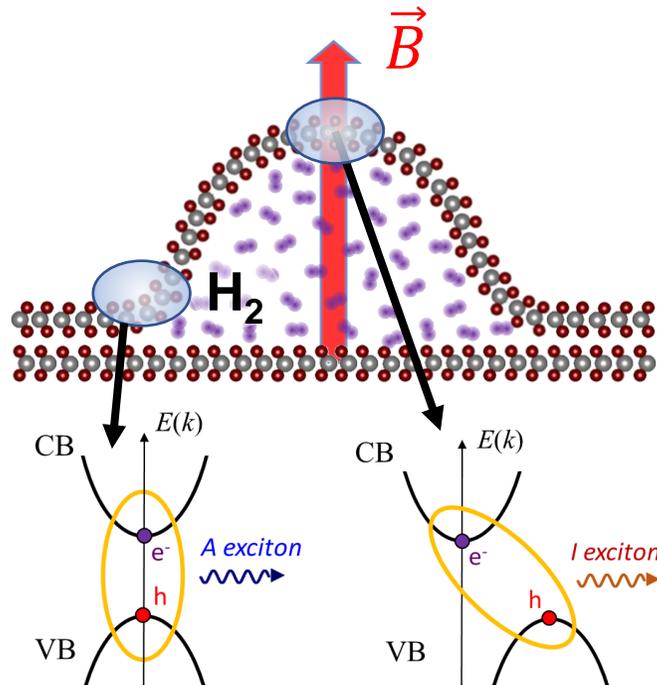


$$\Delta E_Z^{A(I)} = g_{A(I)} \mu_B B$$

$$g_A = 2[L_Z(CB_+) - L_Z(VB_+)]$$

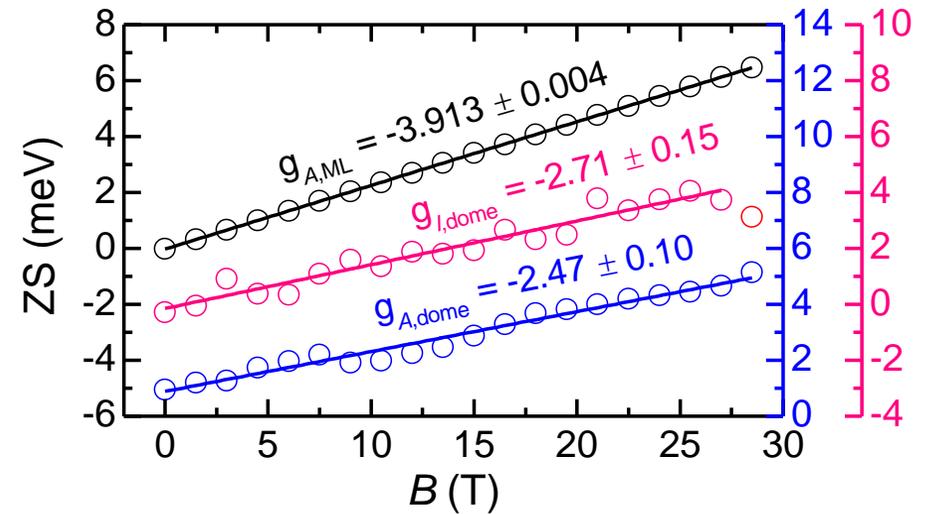
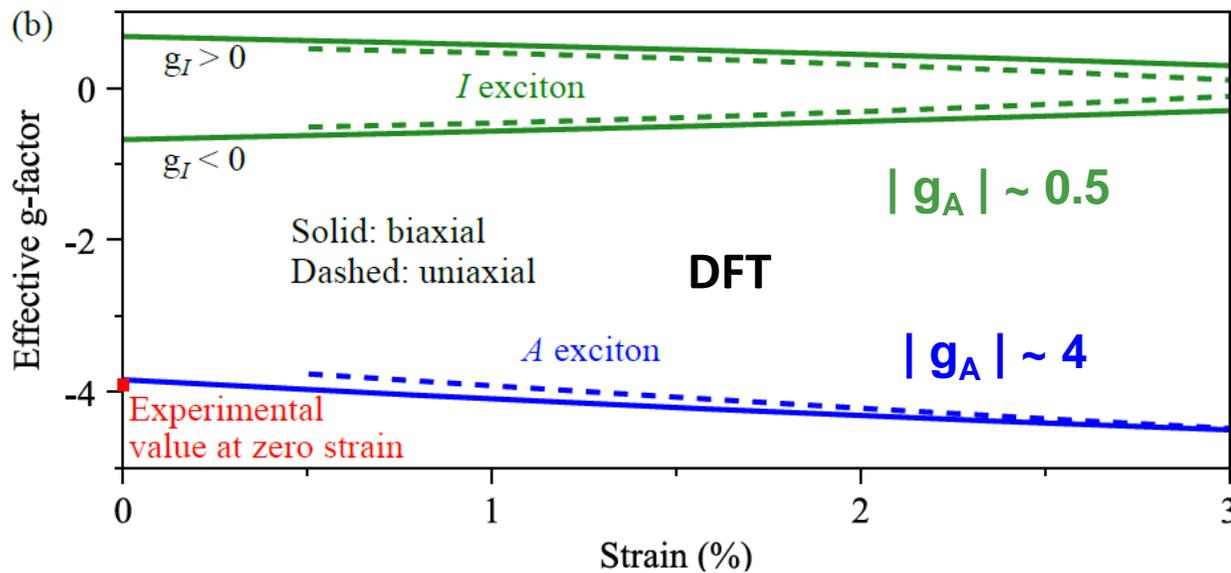
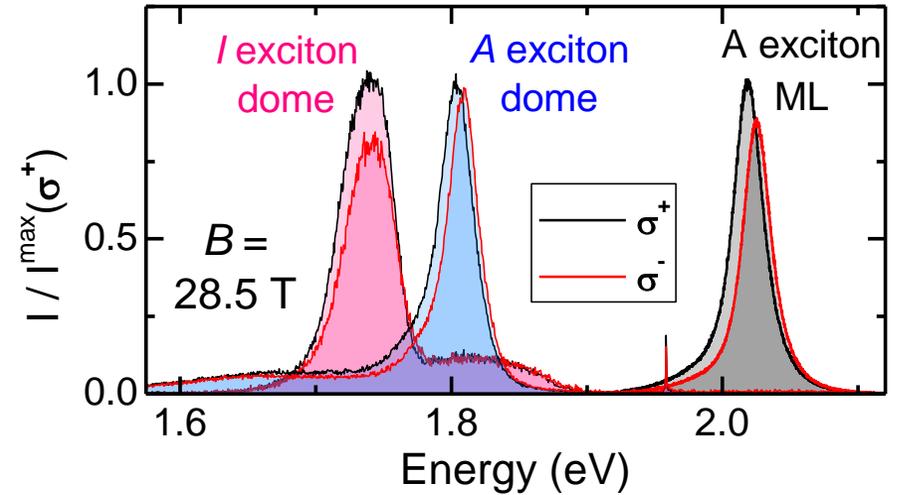
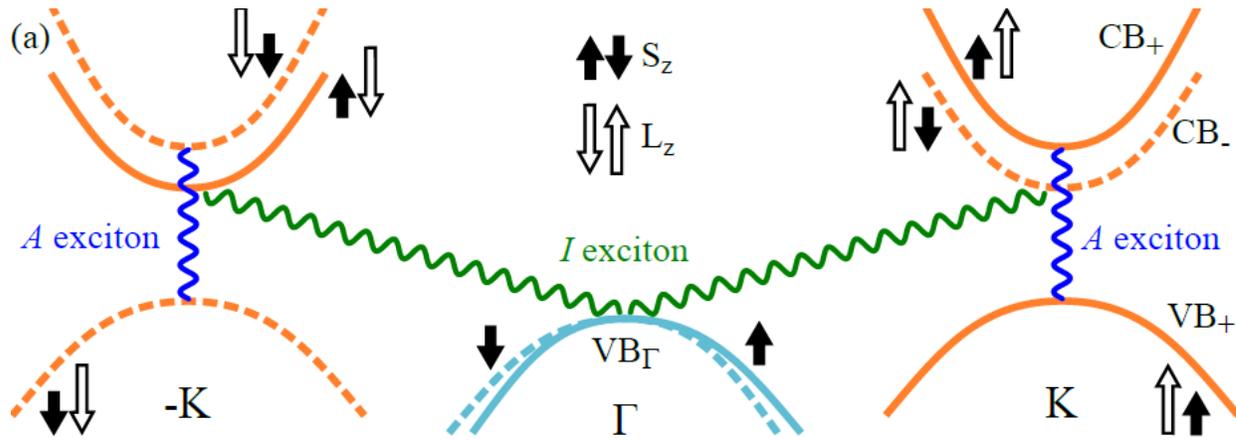
$$g_I = 2[L_Z(CB_-) - 2]$$

How does $g_{A(I)}$ change with strain?



$g_{A(I)}$ should not change with strain
 A and I excitons should show very different splittings

Magneto-PL measurements of WS_2 domes

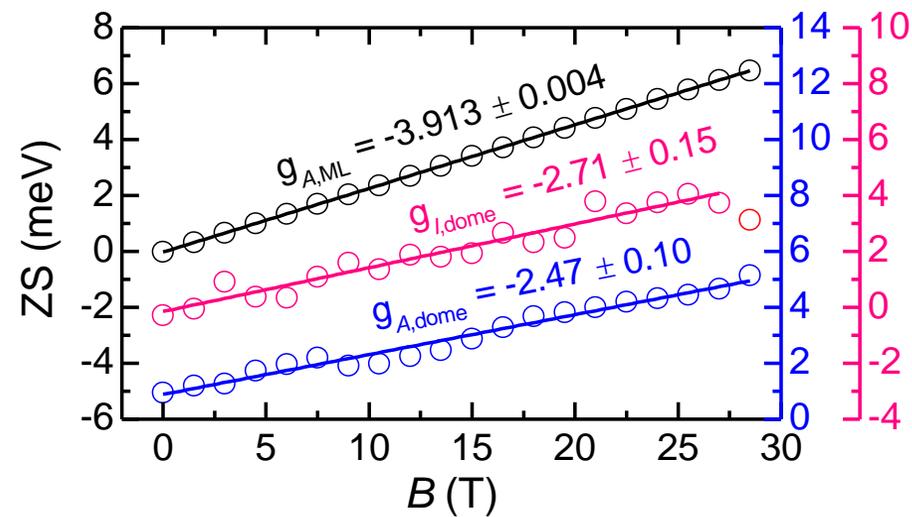
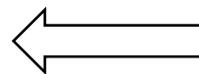
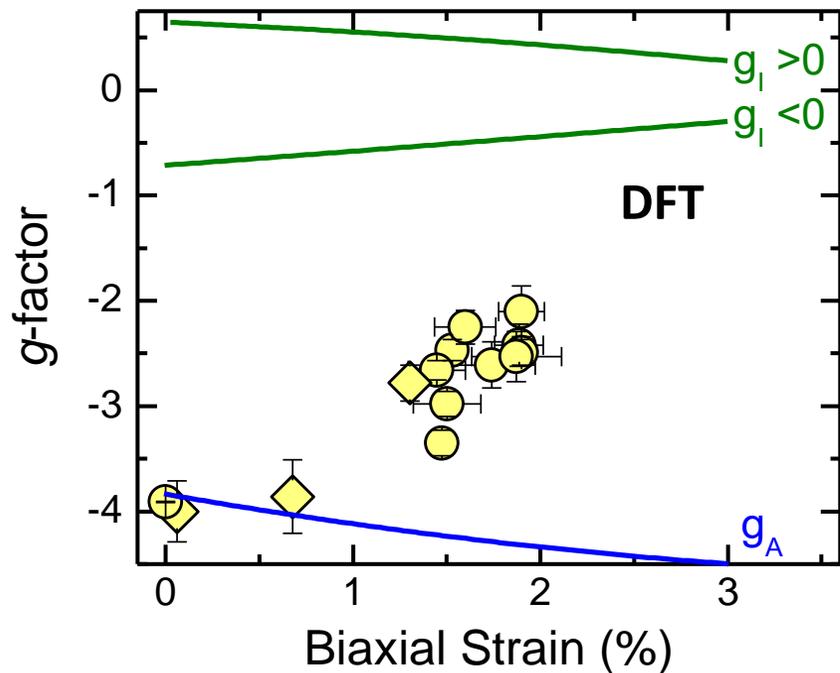
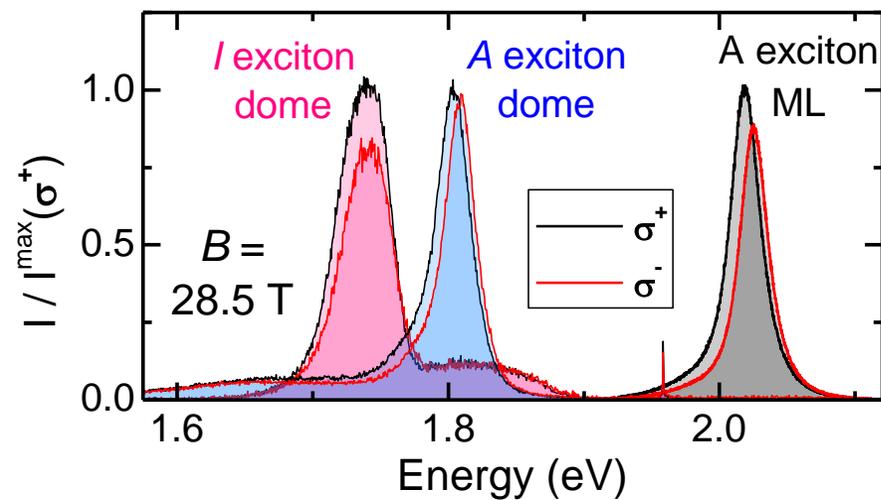
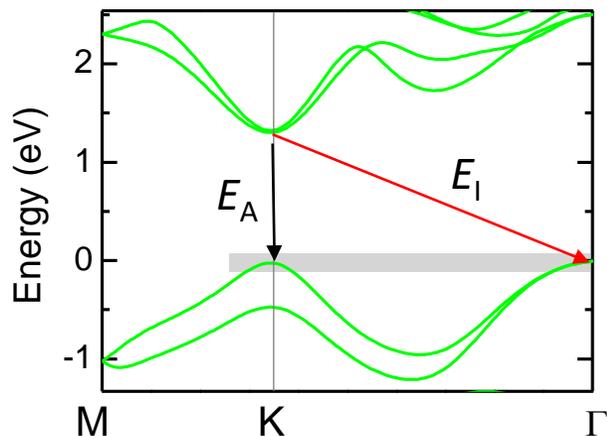


A and I excitons show very similar splittings

$g_{A(I)}$ should not change with strain

Theory-wise A and I excitons should show very different splittings

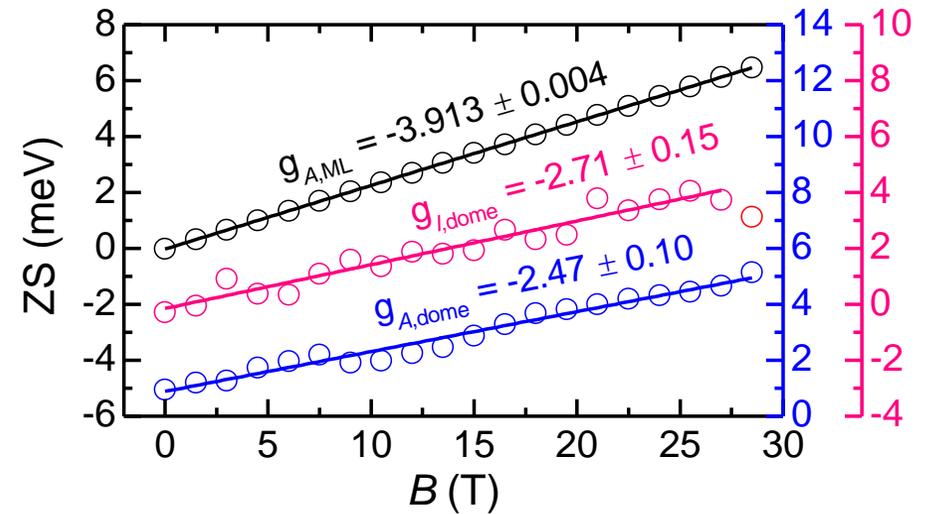
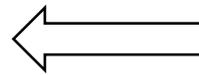
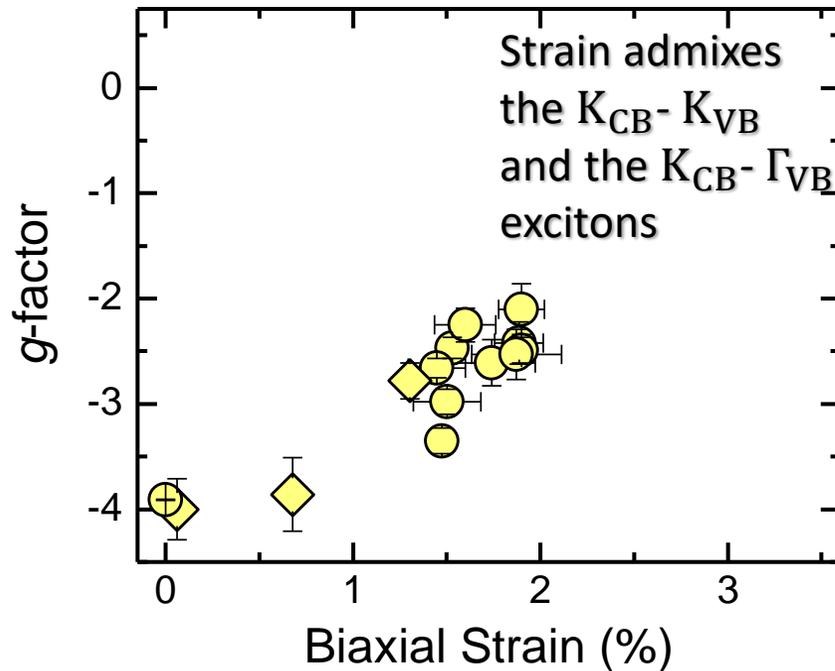
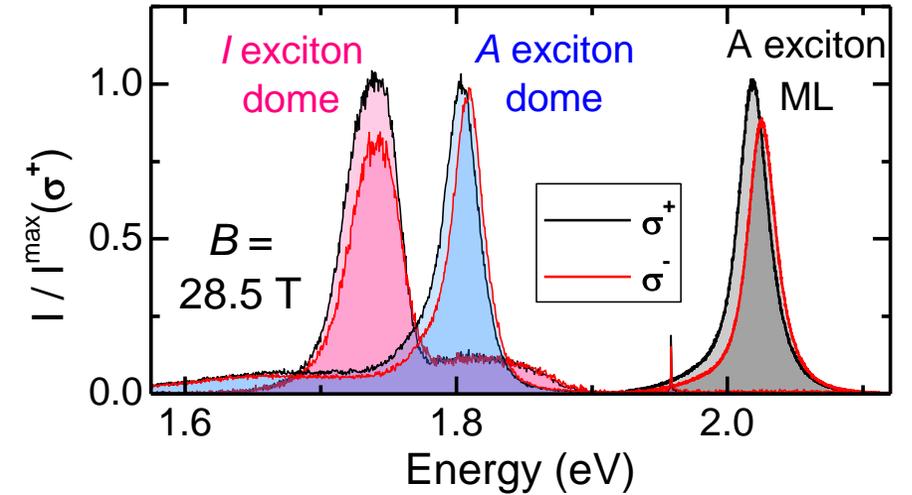
Magneto-PL measurements of WS_2 domes



Magneto-PL measurements of WS_2 domes

Strain-induced A and I exciton hybridisation

$$\begin{bmatrix} E_A + \frac{1}{2}g_A\mu_B B & 0 & \Delta_1 & \Delta_2 \\ 0 & E_A - \frac{1}{2}g_A\mu_B B & -\kappa\Delta_2^* & \kappa\Delta_1^* \\ \Delta_1^* & -\kappa\Delta_2 & E_I + \frac{1}{2}g_I\mu_B B & 0 \\ \Delta_2^* & \kappa\Delta_1 & 0 & E_I - \frac{1}{2}g_I\mu_B B \end{bmatrix}$$

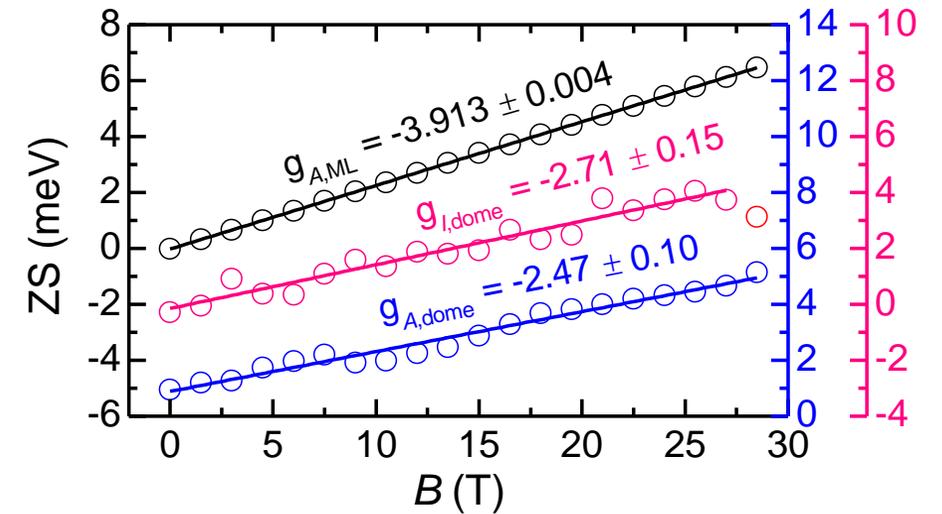
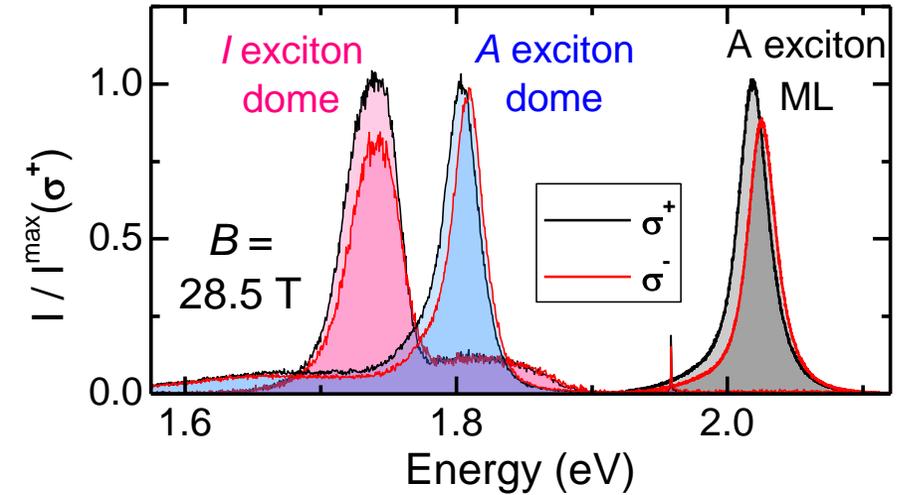
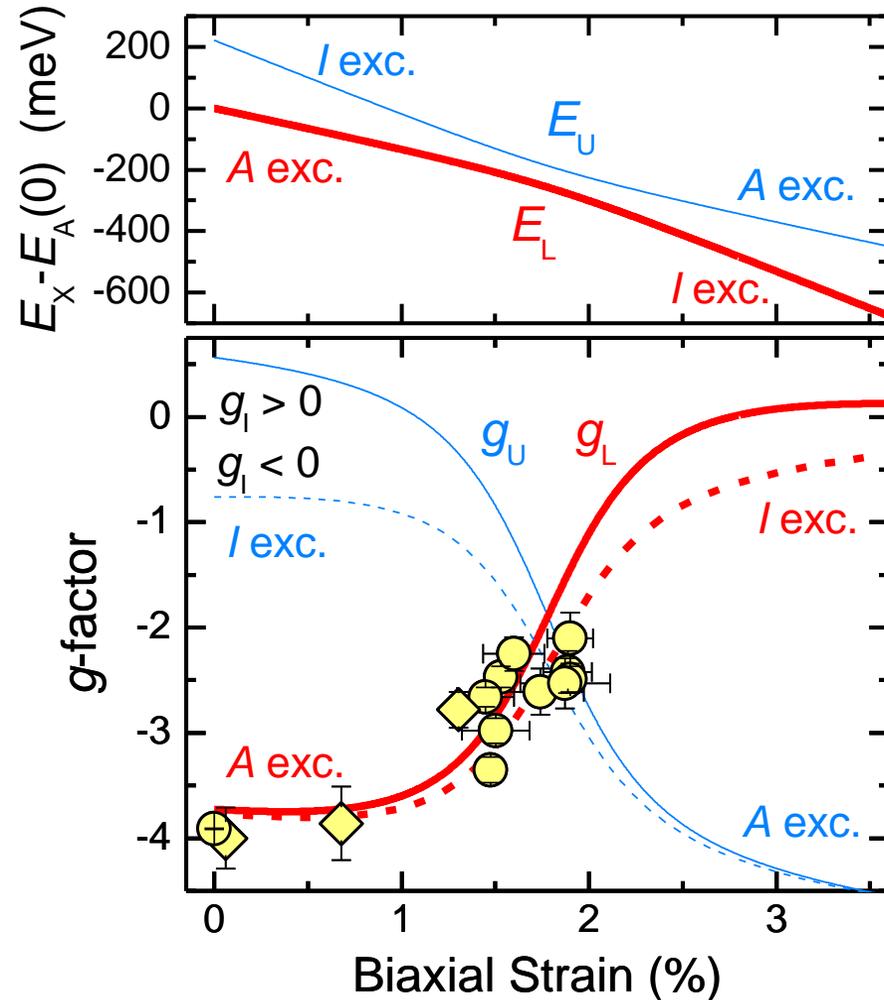


Magneto-PL measurements of WS_2 domes

$$\begin{bmatrix} E_A + \frac{1}{2}g_A\mu_B B & 0 & \Delta_1 & \Delta_2 \\ 0 & E_A - \frac{1}{2}g_A\mu_B B & -\kappa\Delta_2^* & \kappa\Delta_1^* \\ \Delta_1^* & -\kappa\Delta_2 & E_I + \frac{1}{2}g_I\mu_B B & 0 \\ \Delta_2^* & \kappa\Delta_1 & 0 & E_I - \frac{1}{2}g_I\mu_B B \end{bmatrix}$$

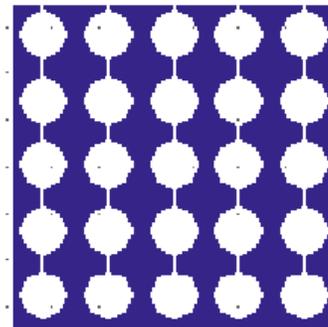
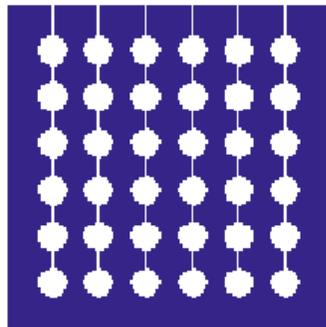
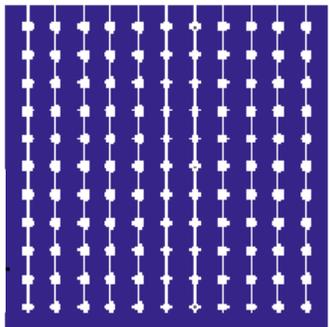
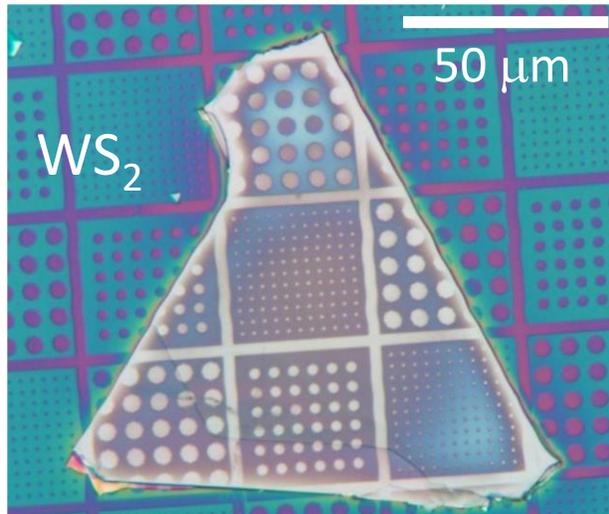
$$\Delta_1 = 35 \text{ meV}$$

$$\Delta_2 = 0$$



Zeeman-splitting measurements unveil exciton hybridization
 Phonons? Disorder?

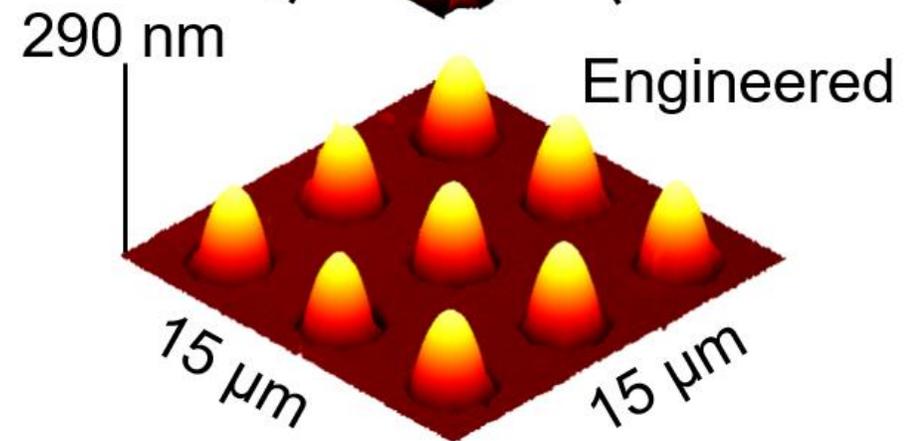
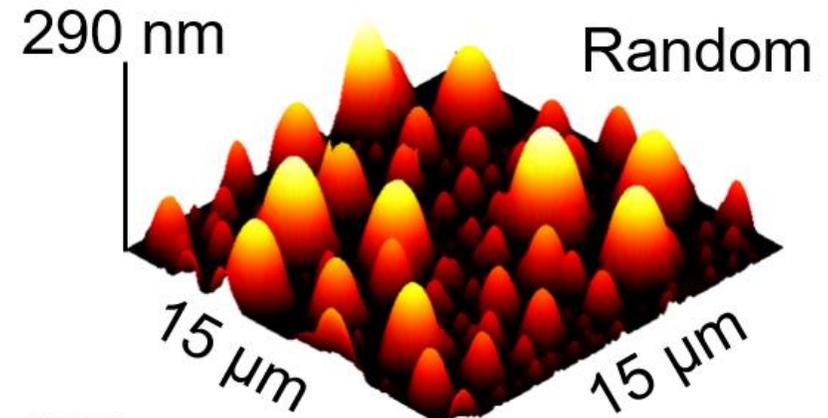
Controlled dome formation



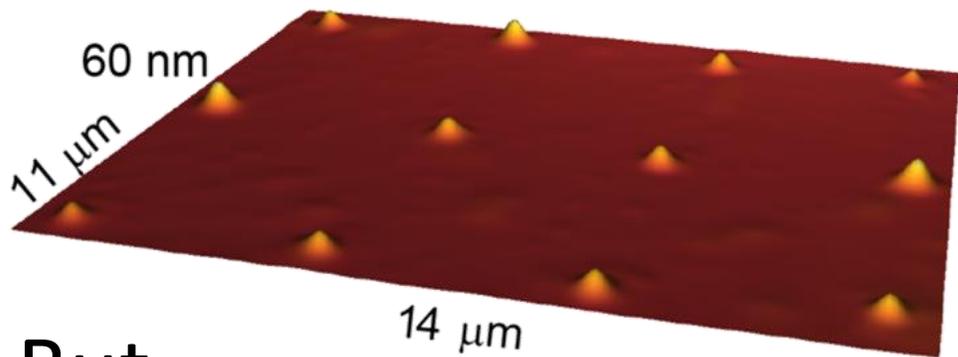
$S = 1 \mu\text{m}$

$3 \mu\text{m}$

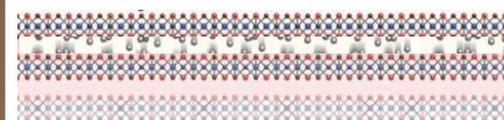
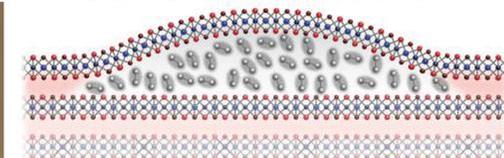
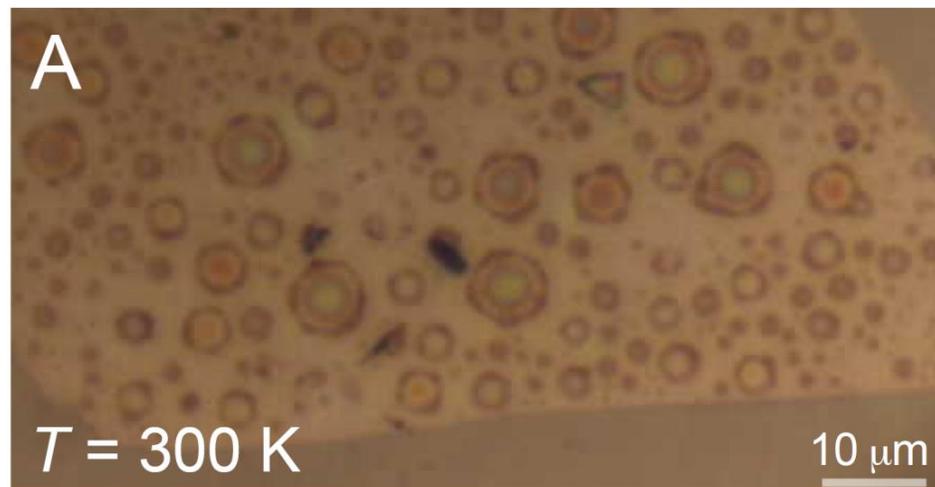
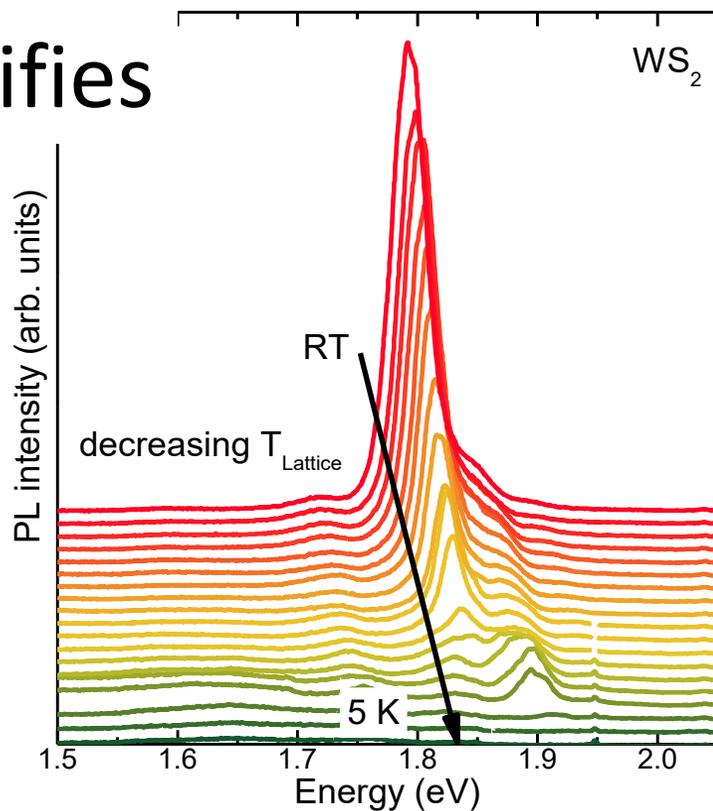
$5 \mu\text{m}$



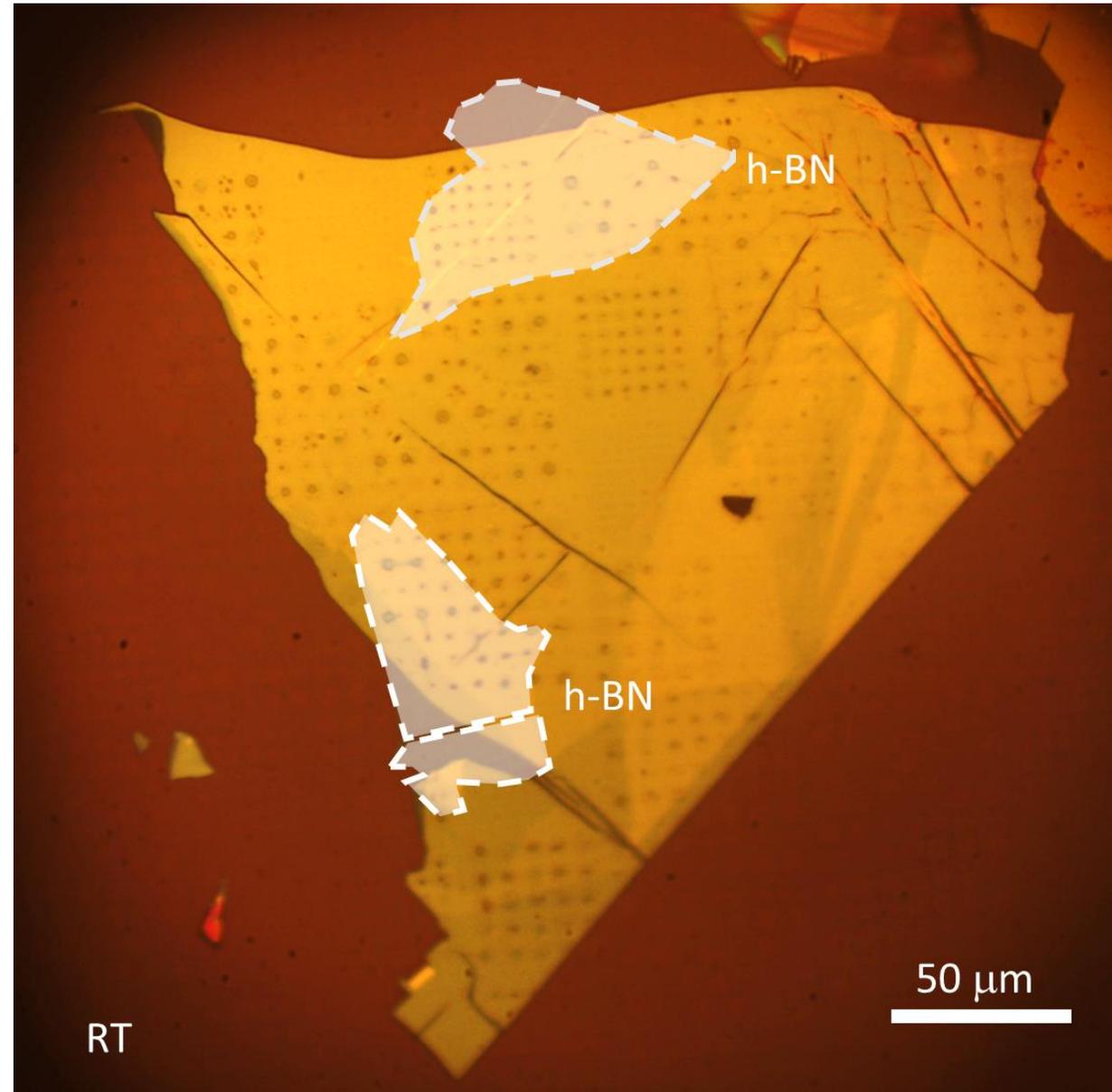
Space-controlled emitters?



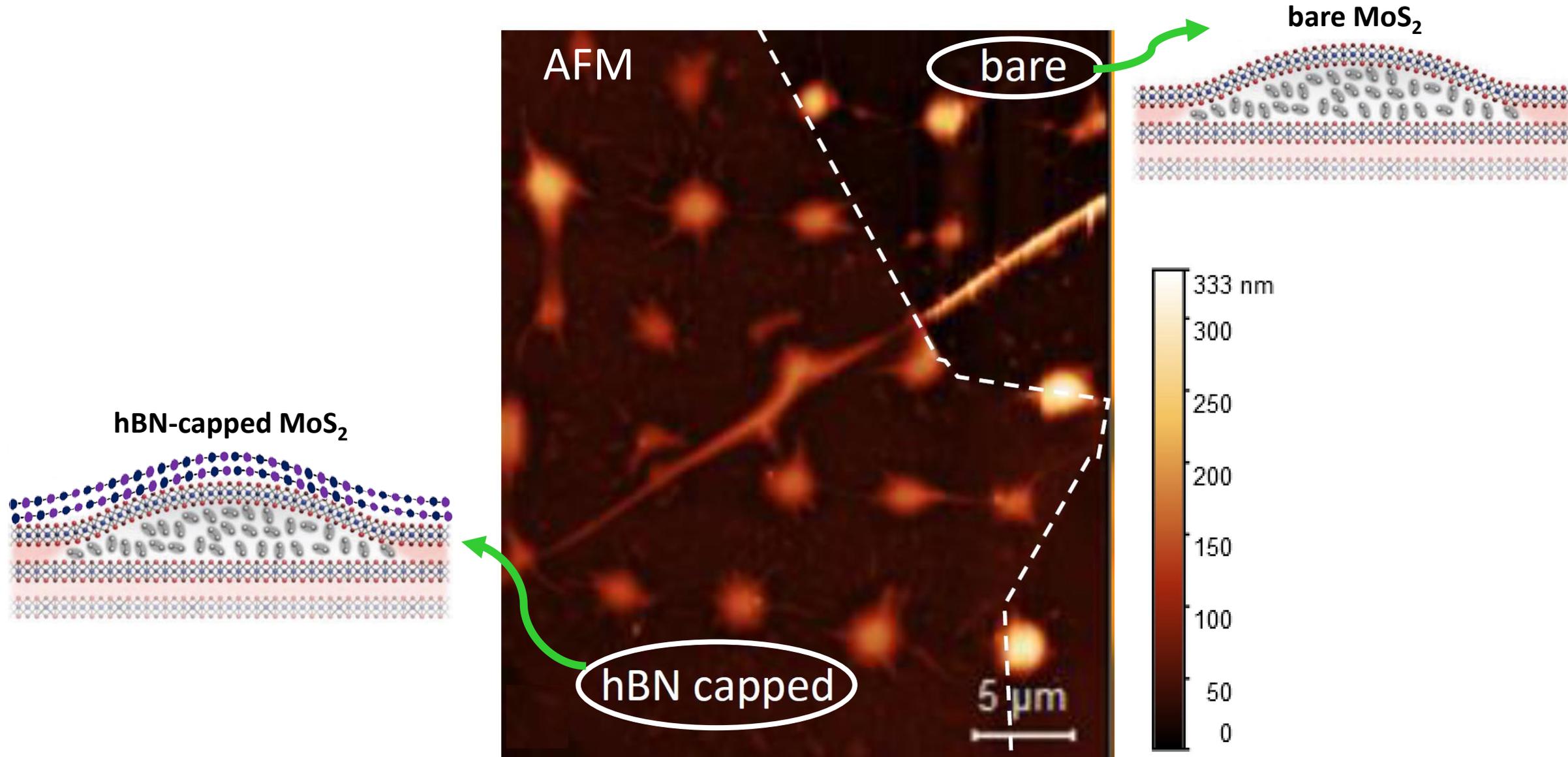
But...
 H_2 liquifies



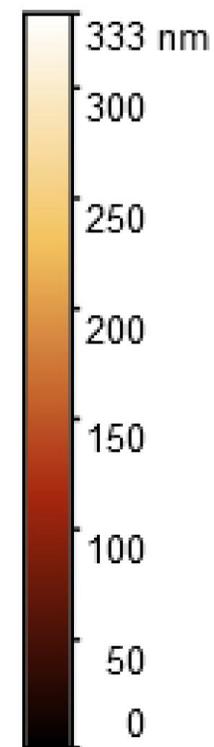
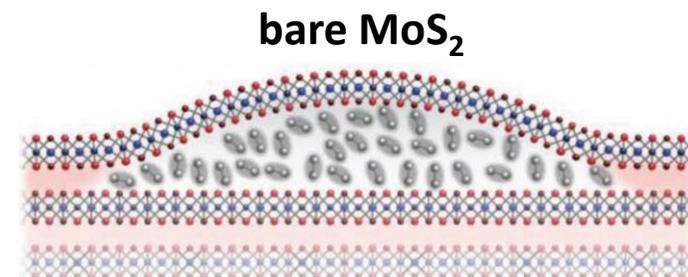
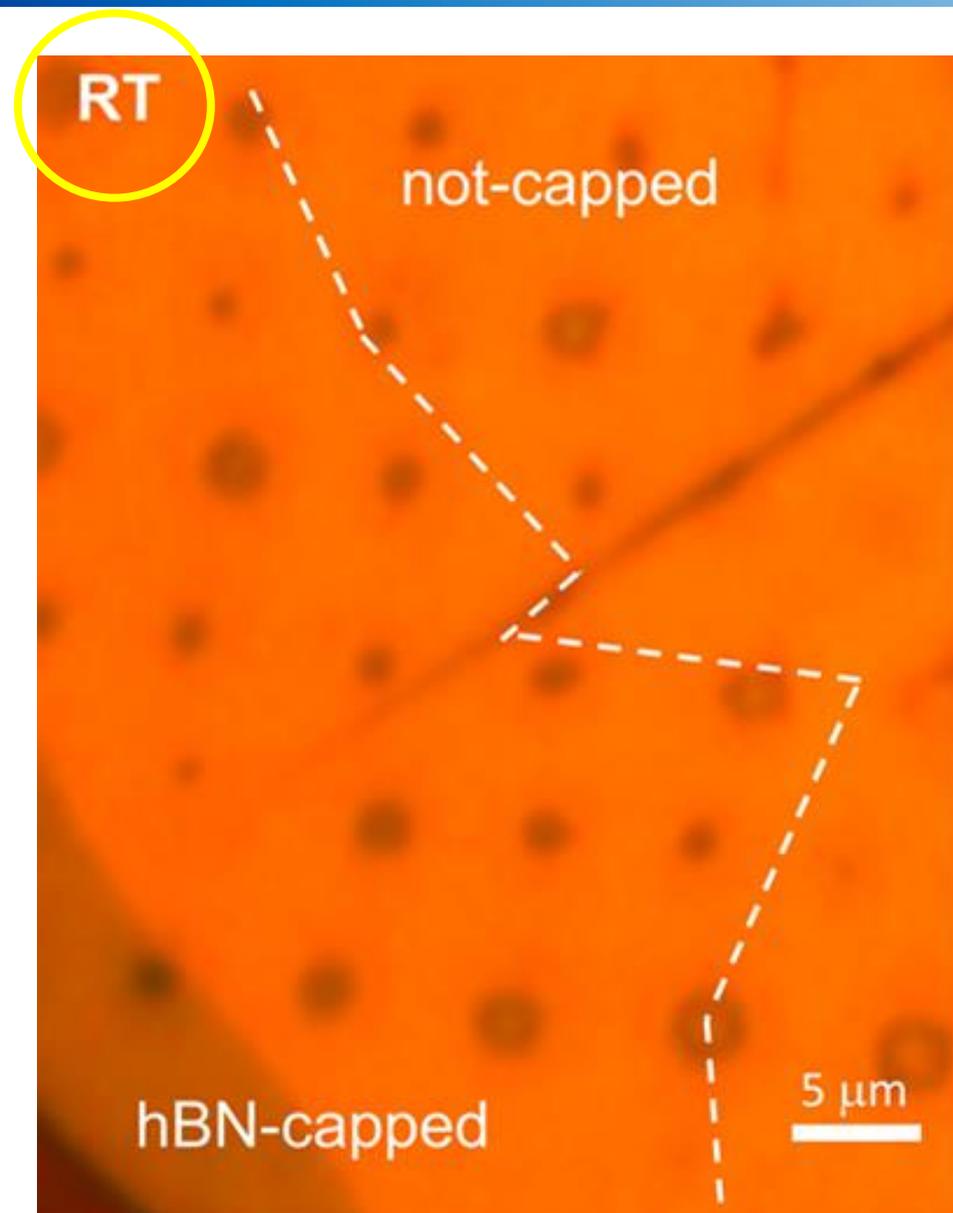
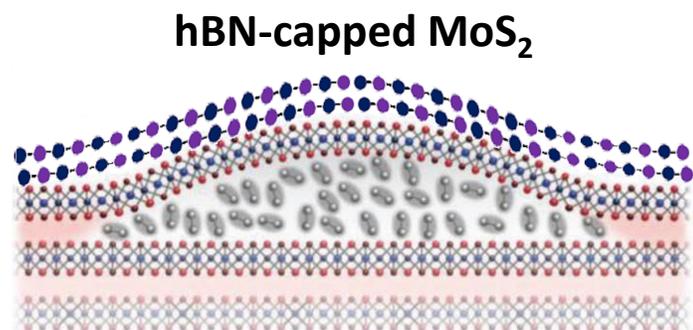
How to circumvent H_2
liquefaction?



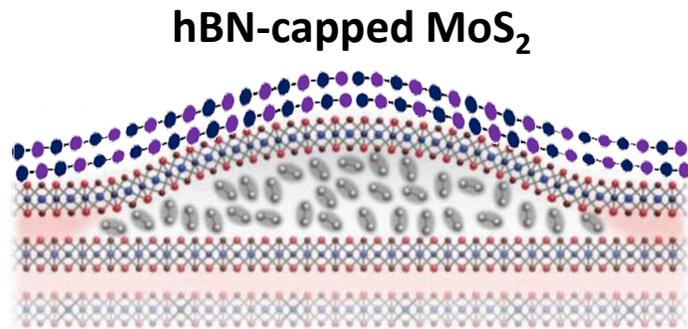
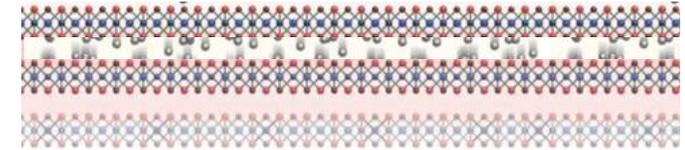
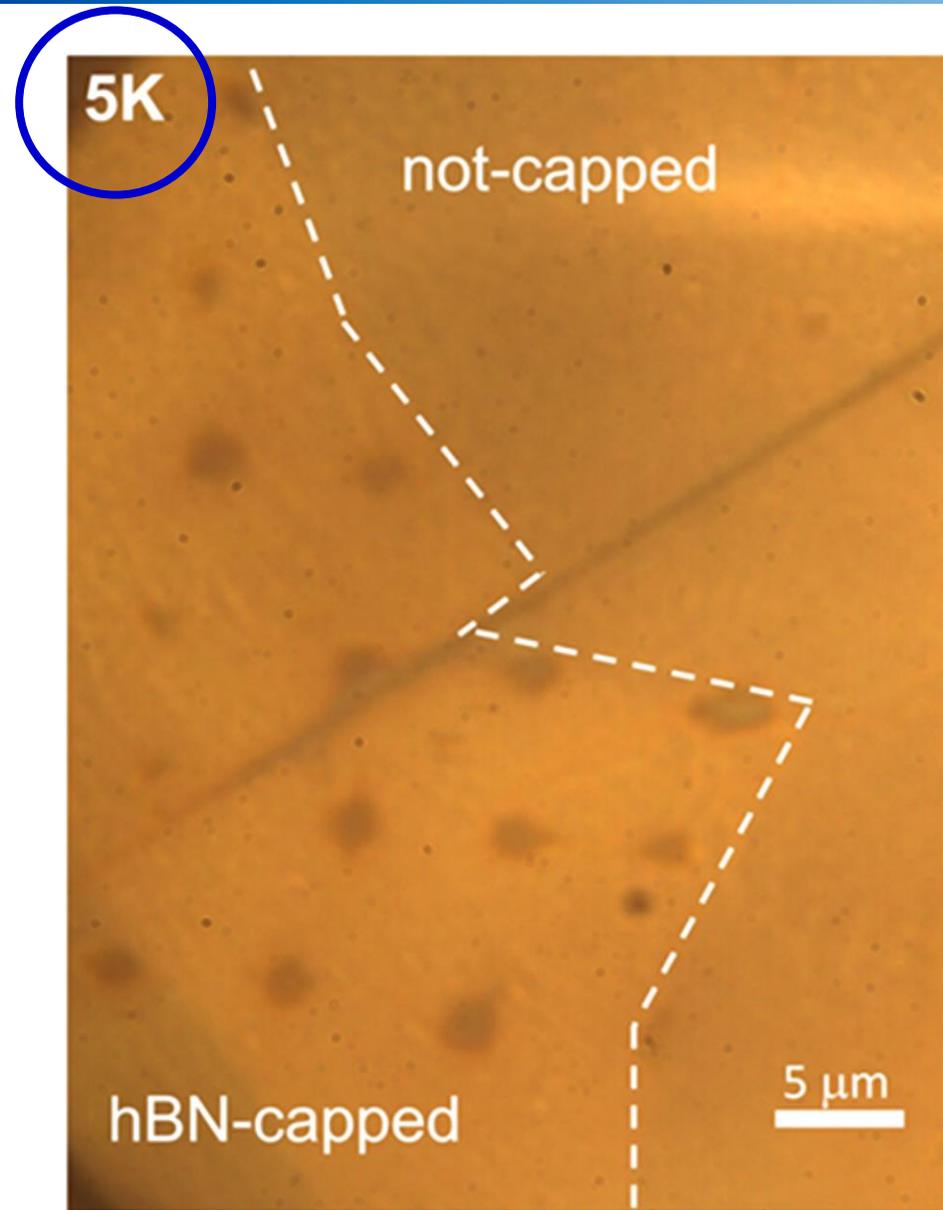
hBN heterostructuring



hBN heterostructuring



hBN heterostructuring

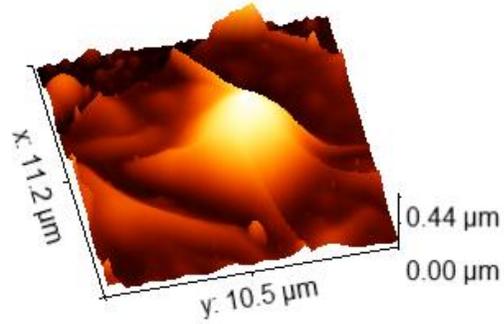
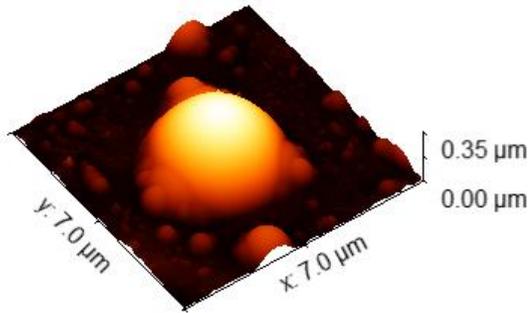


capped domes do not deflate

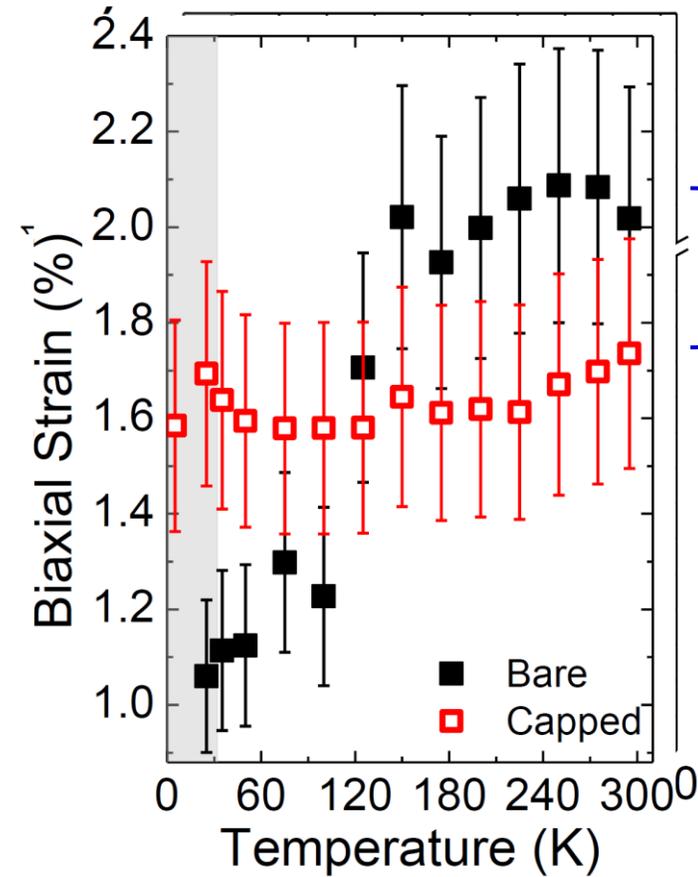
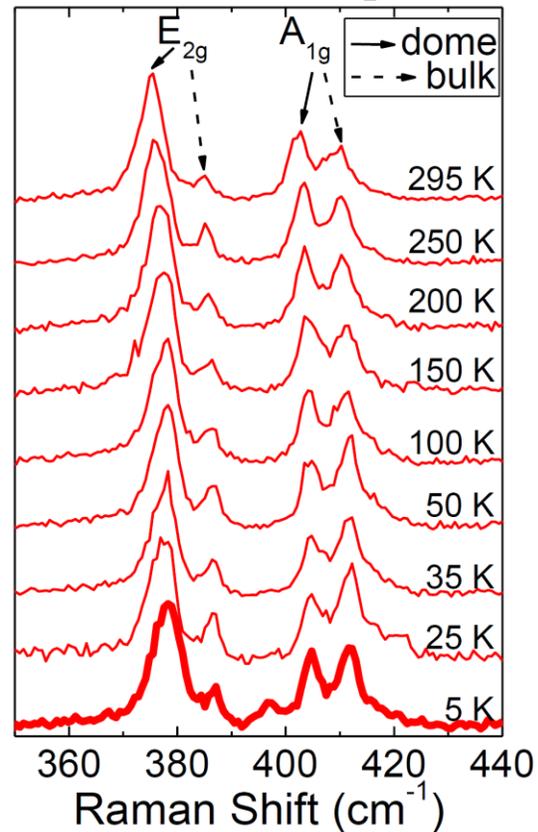
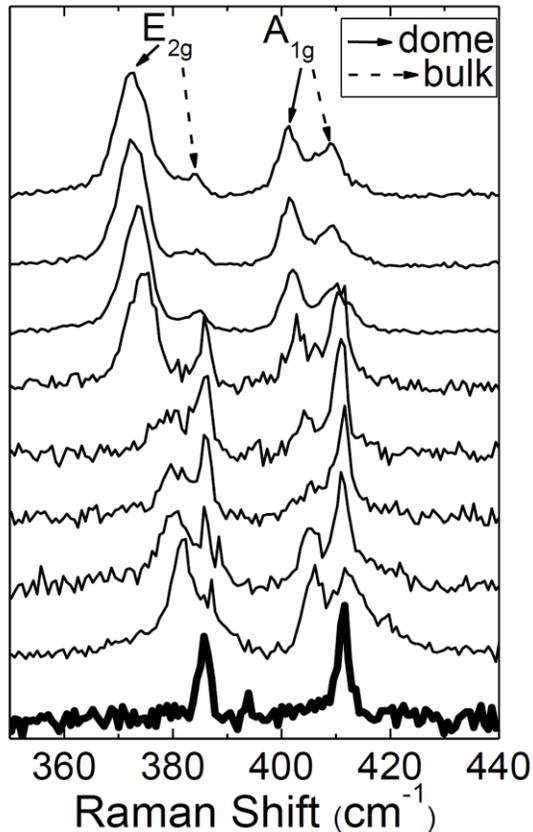
hBN heterostructuring

bare MoS₂ dome

hBN-capped MoS₂ dome



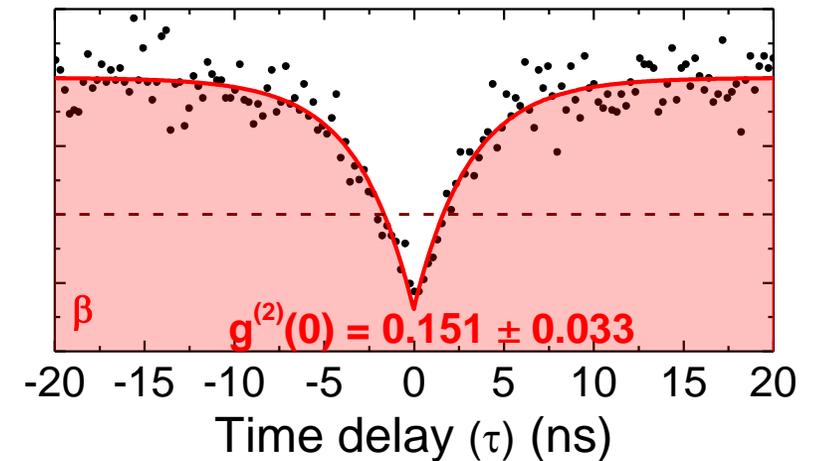
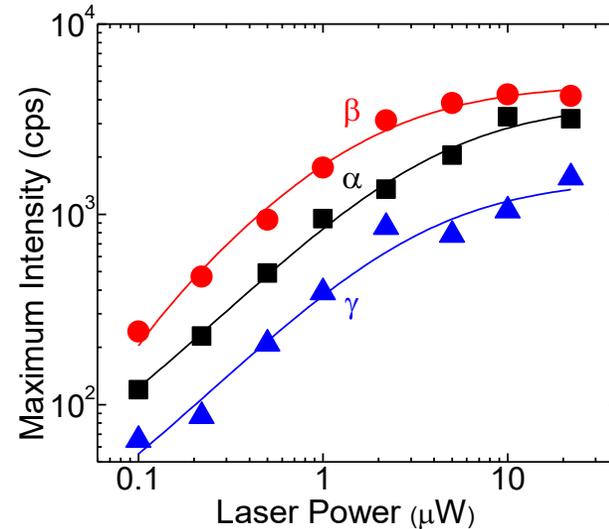
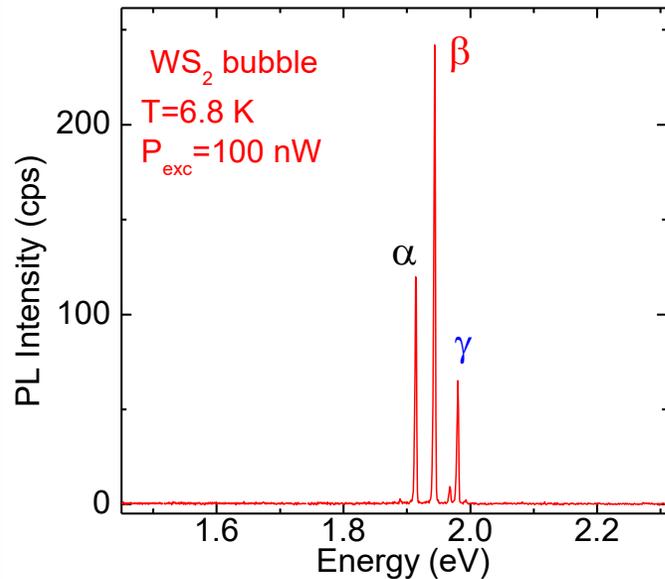
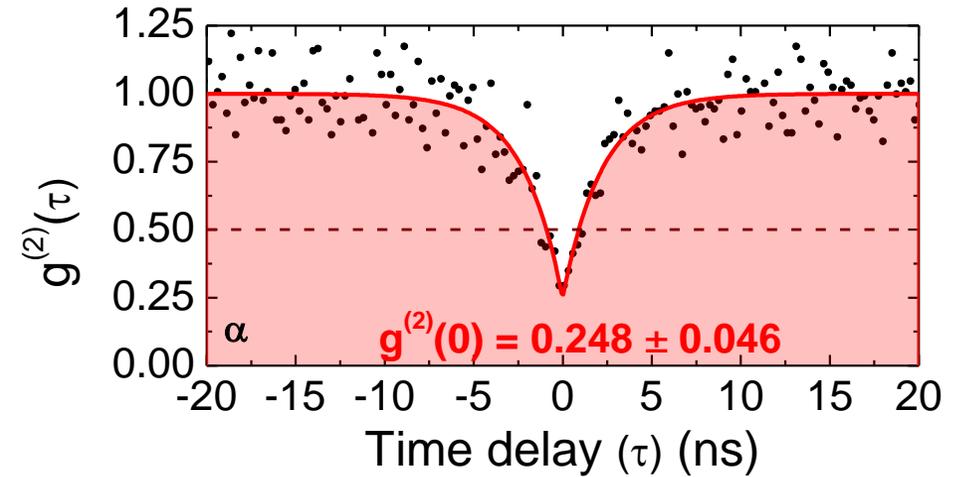
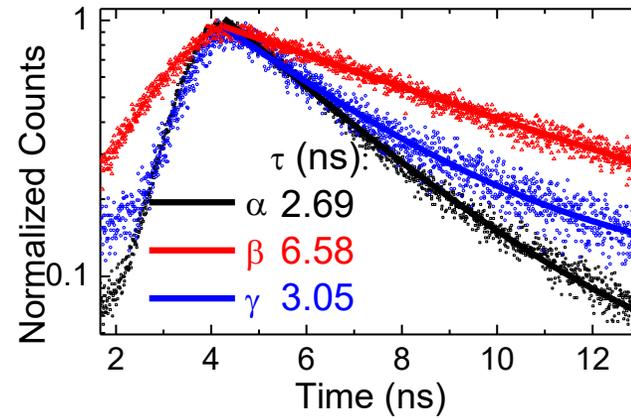
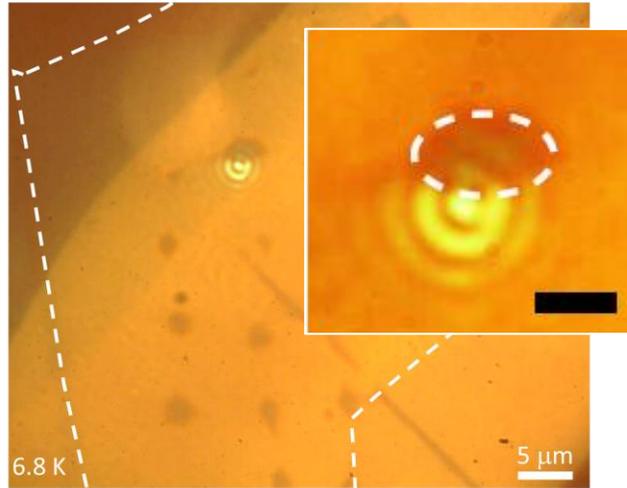
Normalized Raman Intensity (arb. units)



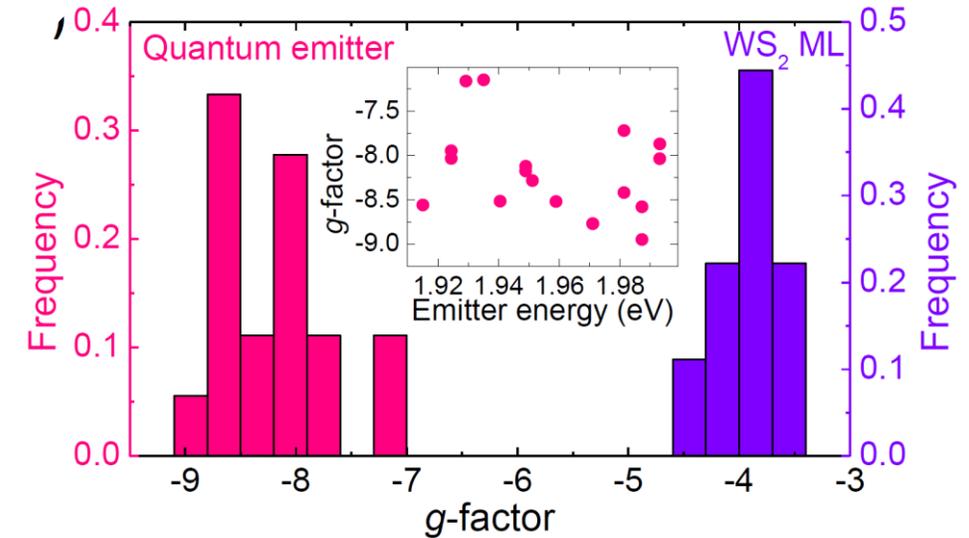
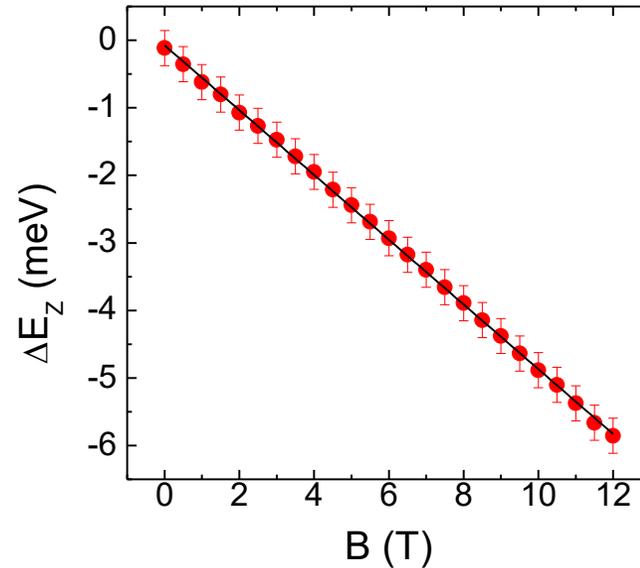
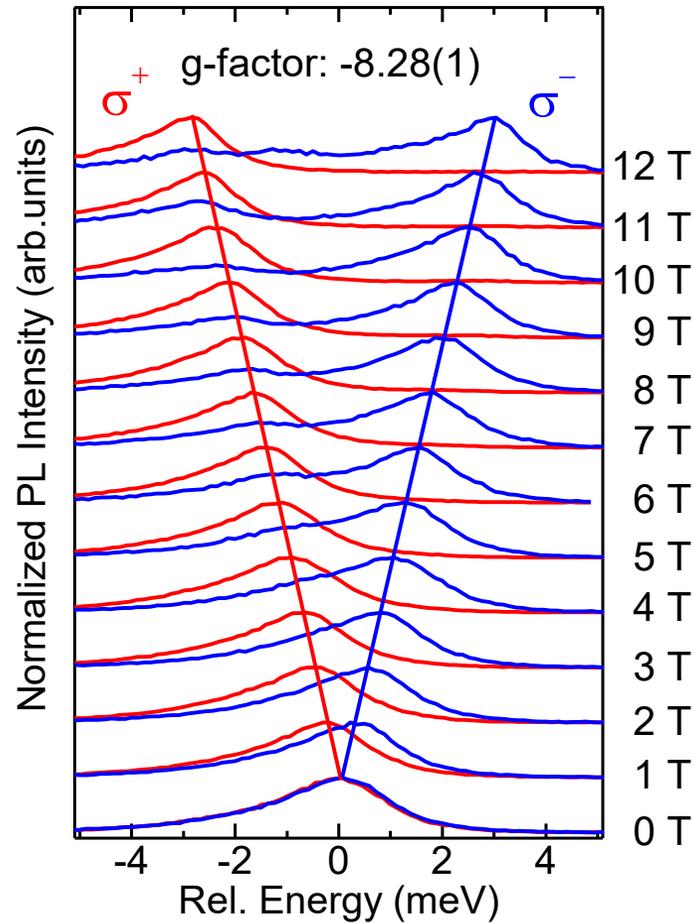
hBN capping causes strain decrease: Elastic energy transfer

$$\varepsilon(T) = \frac{1}{2} \cdot \left(\frac{\omega_{E_{2g}}^{ML}(T) - \omega_{E_{2g}}^{dome}(T)}{\omega_{E_{2g}}^{ML}(T)} \right) \frac{1}{\gamma_{E_{2g}}}$$

Space-controlled quantum emitters



Space-controlled quantum emitters

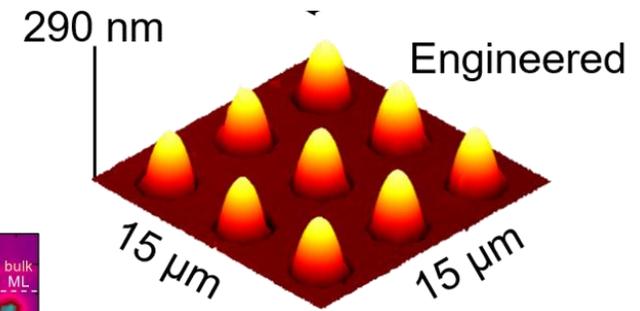


Like for WSe₂ quantum emitters, the large g value indicates the involvements of an electron in a defect state and a hole in the valence band

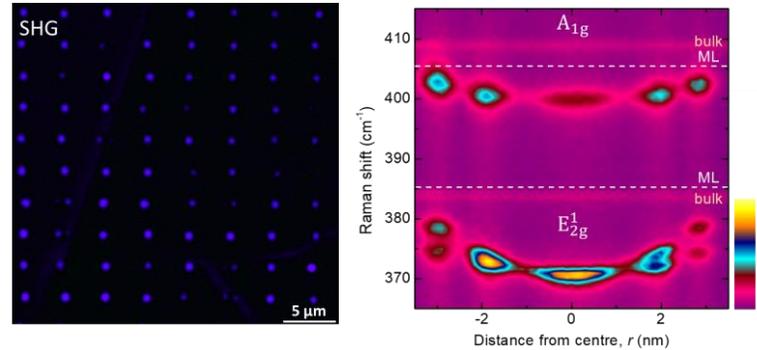
$$\Delta E_Z = g_{\text{exc}} \mu_B B$$

Conclusions

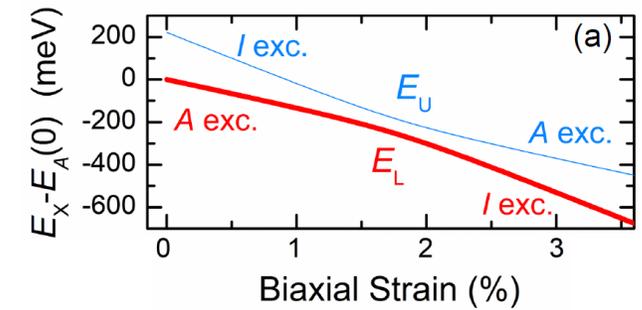
- Durable, spatially controlled domes can be created in TMDs and hBN



- The domes host complex strain field and act as efficient light emitters



- Complex strain fields give access to exciton hybridization phenomena



- Applications for site-controlled quantum light sources

